Full Length Research Paper

# Construction of some new exact structures for the nonlinear lattice equation 

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Received 23 October, 2014; Accepted 13 November, 2014


#### Abstract

In the present work we examine a generalized coth and csch functions method to construct new exact travelling solutions to the nonlinear lattice equation. The technique of the homogeneous balance method is used to handle the appropriated solutions. Some exact solutions obtained are new.


Key words: Nonlinear lattice, travelling wave solutions.

## INTRODUCTION

Several methods have been developed for analytic solving of nonlinear partial differential equations. Specially, almost all of these nonlinear model equations were appeared (Wang and Li, 2008; Korteweg and Vries, 1995; Khelil et al., 2006) to give different structures to the solutions. Besides traditional methods such as autoBacklund transformation, Lie Groups, inverse scattering transformation and Miura's transformation, a vast variety of the direct methods for obtaining explicit travelling solitary wave solutions have been found (Zerarka and Foester, 2005; Ibrahim and El-Kalaawy, 2007; Lü, $2014 \mathrm{a}, \mathrm{b}$ ). The availability of symbolic computation packages can be facilitating many direct approaches to establish solutions to non-linear wave equations (Xu and Zhang, 2007; Özis and Yıldırım, 2008). Various extension forms of the sine-cosine and tanh methods proposed by Malfliet and Wazwaz have been applied to solve a large class of nonlinear equations (Malfliet 1996a,b; Wazwaz 2004; He and Wu, 2006a,b). More importantly, another mathematical treatment is established and used in the
analysis of these nonlinear problems, such as Jacobian elliptic function expansion method, the variational iteration method, pseudo spectral method, the averaging method, and many others powerful methods (Odibat and Momani, 2006; Rafei and Ganji, 2006; Yu, 2007; Zhu, 2007a,b; Lü and Peng, 2013a,b,c; Lü, 2013; Lü et al., 2010; Jia et al., 2014; Liu and Qian, 2011). The aim of this work is to propose an efficient approach to examine new developments in a direct manner without requiring any additional condition on the investigation of exact solutions with the coth and csch functions method for a lattice system. We expect that the presented method could lead to construct successfully many other solutions for a large variety of other nonlinear evolution equations.

## ANALYSIS OF THE PROBLEM

We consider the following nonlinear problem for the lattice equation as:

[^0]$R\left(u, u_{t}, u_{x}, u_{y}, u_{z}, u_{x y}, u_{y z}, u_{x z}, u_{x x}, \ldots\right)=0$,
Here the subscripts represent partial derivatives, and $u(t, x, y, z, \ldots)$ is an unknown function to be determined. We take the following transformation for the new wave variable as:
$\xi=\sum_{i=0}^{p} \alpha_{i} \chi_{i}+\delta$,
$\chi_{i}$ are distinct variables, and when $p=1$, $\xi=\alpha_{0} \chi_{0}+\alpha_{1} \chi_{1}+\delta$, the quantities $\alpha_{0}, \alpha_{1}$ are called the wave pulsation $\omega$ and the wave number $k$ respectively if $\chi_{0}, \chi_{1}$ are the variables $t$ and $x$ respectively. In the discrete case for the position $x$ and with continuous variable for the time $t, \xi$ becomes with some modifications $\xi_{n}=n d+c t+\delta$ and $n$ is the discrete variable. $d$ and $\delta$ are arbitrary constants and $c$ is the velocity. We use the traveling wave reduction transformation for Equation (1) as:
$u\left(\chi_{0}, \chi_{1}, \ldots\right)=U(\xi)$,
and the chain rule
\[

$$
\begin{equation*}
\frac{\partial}{\partial \chi_{i}}(.)=\alpha_{i} \frac{d}{d \xi}(.), \frac{\partial^{2}}{\partial \chi_{i} \partial \chi_{j}}(.)=\alpha_{i} \alpha_{j} \frac{d^{2}}{d \xi^{2}}(.), \cdots, \tag{4}
\end{equation*}
$$

\]

Upon using Equations (3) and (4), the nonlinear problem (1) becomes an ODE like

$$
\begin{equation*}
Q\left(U, U_{\xi}, U_{\xi \xi}, U_{\xi \xi \xi}, U_{\xi \xi \xi \xi}, \ldots\right)=0 \tag{5}
\end{equation*}
$$

## APPLICATIONS

The one-dimensional lattice equation (Zhu, 2007, 2008) is written as:
$\frac{d u(n, t)}{d t}-\left(a+b u(n, t)+u^{2}(n, t)\right)[u(n+1, t)-u(n-1, t)]=0$,
We first combine the independent variables, into a wave variable using $\xi_{n}$ as

$$
\begin{equation*}
\xi_{n}=n d+c t+\delta \tag{7}
\end{equation*}
$$

and we take the travelling wave solutions of the system
(6) using Equation (7) as $u(n, t)=U\left(\xi_{n}\right)$. By using the chain rule (4), the system (6) can be obtained as follows:
$c U_{\xi_{n}}\left(\xi_{n}\right)-\left(a+b U\left(\xi_{n}\right)+U^{2}\left(\xi_{n}\right)\right)\left(U\left(\xi_{n+1}\right)-U\left(\xi_{n-1}\right)\right)=0$,
Where subscript denotes the differential with respect to $\xi_{n}$.

## THE COTH FUNCTION METHOD

Suppose that Equation (8) has the following solution:
$U\left(\xi_{n}\right)=\sum_{j=0}^{M} A_{j} \operatorname{coth}^{j}\left(\xi_{n}\right)$,
Where $M$ is an undetermined integer and $A_{j}$ are coefficients to be determined later. In order to determine values of the parameter $M$, we balance the linear term of highest order in Equation (8) with the highest order nonlinear term. By simple calculation, we have $2 M=M+1$ and the solution (9) takes the form
$U\left(\xi_{n}\right)=A_{0}+A_{1} \operatorname{coth}\left(\xi_{n}\right)$,
Substituting the solution (10) into Equation (8), and equating to zero the coefficients of all powers of $\operatorname{coth}^{j}\left(\xi_{n}\right)$ yields a set of algebraic equations for $A_{0}$, $A_{1}$ and $c$ as:

$$
\begin{align*}
b A_{1}+2 A_{0} A_{1} & =0 \\
2\left(a+b A_{0}+A_{0}^{2}\right) & =c \operatorname{coth}(d)  \tag{11}\\
-2 A_{1}^{2} \operatorname{coth}(d) & =c
\end{align*}
$$

Solving the system of algebraic equations with the aid of Mathematical, we obtain

$$
\begin{align*}
A_{0} & =-\frac{b}{2}, \\
A_{1} & = \pm \frac{\tanh (d)}{2} \sqrt{b^{2}-4 a},  \tag{12}\\
c & =\frac{\tanh (d)}{2}\left(4 a-b^{2}\right),
\end{align*}
$$

and the two travelling wave solutions of the problem of interest follow

$$
\begin{equation*}
U_{ \pm}\left(\xi_{n}\right)=-\frac{b}{2} \pm \frac{\tanh (d)}{2} \sqrt{b^{2}-4 a} \operatorname{coth}\left[n d+\frac{\tanh (d)}{2}\left(4 a-b^{2}\right) t+\delta\right] \tag{13}
\end{equation*}
$$



Figure 1. The graphs show the wave solutions of $u_{ \pm}(n, t)=U_{ \pm}\left(\xi_{n}\right)$ in Equation (13): (a) solution $u_{+}(n, t)=U_{+}\left(\xi_{n}\right)$,
(b) solution $u_{-}(n, t)=U_{-}\left(\xi_{n}\right)$. For both curves: $a=2, b=3, d=1, \delta=0$.

Where $d$ and $\delta$ are arbitrary constants.
Figure 1(a) and (b) show the physical waves $u_{+}(n, t)=U_{+}\left(\xi_{n}\right)$ and $u_{-}(n, t)=U_{-}\left(\xi_{n}\right)$ in Equations (13).

## THE COTH-CSCH FUNCTION METHOD

The solutions of Equation (8) can be expressed in the form
$U\left(\xi_{n}\right)=\alpha+\sum_{j=1}^{M} A_{j} \operatorname{coth}^{j}\left(\xi_{n}\right)+B_{j} \operatorname{csch}^{j}\left(\xi_{n}\right)$,
Where $\alpha, A_{j}$ and $B_{j}$ are parameters to be determined. The parameter $M$ is found by balancing the highestorder linear term with the nonlinear terms, we obtain $M=1$, and $U\left(\xi_{n}\right)$ becomes
$U\left(\xi_{n}\right)=\alpha+A \operatorname{coth}\left(\xi_{n}\right)+B \operatorname{csch}\left(\xi_{n}\right)$,
Substituting Equation (15) into the relevant nonlinear differential Equation (8) and with the help of Mathematical we get a system of algebraic equations with respect to $c$, $\alpha, A$ and $B$.

$$
\begin{align*}
& A c+2 A\left(A^{2}+B^{2}\right) \operatorname{coth}(d)+4 A B^{2} \operatorname{csch}(d)=0, \\
& 4 A^{2} B \operatorname{coth}(d)+B c \operatorname{coth}(d)^{2}-2 B\left(\alpha^{2}-B^{2}+a+\alpha b\right) \operatorname{csch}(d)=0, \\
&-(B c)-4 A^{2} B \operatorname{coth}(d)-2 B\left(A^{2}+B^{2}\right) \operatorname{csch}(d)=0, \\
& 2 A(2 \alpha B+B b)=0, \\
& 2 A(2 A \alpha+A b) \operatorname{coth}(d)+2 B(2 \alpha B+B b) \operatorname{csch}(d)=0, \\
& 2 A\left(\alpha^{2}-B^{2}+a+\alpha b\right)-A c \operatorname{coth}(d)=0 \tag{16}
\end{align*}
$$

After some algebra, and with the help of Mathematical, the following values for the parameters $c, \alpha, A$, and $B$ are obtained:

## First set

$c=\frac{1}{2}\left(4 a-b^{2}\right) \tanh (d)$,
$\alpha=-\frac{b}{2}$,
$A= \pm \frac{1}{2} \sqrt{b^{2}-4 a} \tanh (d)$,
$B=0$,
and the travelling solutions of Equation (17) are obtained as:
$U_{ \pm}\left(\xi_{n}\right)=-\frac{b}{2} \pm \frac{\tanh (d)}{2} \sqrt{b^{2}-4 a} \operatorname{coth}\left(\xi_{n}\right)$,
Where $\xi_{n}=n d+\frac{\tanh (d)}{2}\left(4 a-b^{2}\right) t+\delta$. The solutions are similar to those obtained by the coth-function method Equation (13).

## Second set

$c=\frac{1}{2}\left(4 a-b^{2}\right) \sinh (d)$,
$\alpha=-\frac{b}{2}$,
$A=0$,
$B= \pm \frac{1}{2} \sqrt{b^{2}-4 a} \sinh (d)$,


Figure 2. The graphs show the wave solution of $U_{ \pm}\left(\xi_{n}\right)$ in (20): (a) solution $u_{+}(n, t)=U_{+}\left(\xi_{n}\right)$, (b) solution $u_{-}(n, t)=U_{-}\left(\xi_{n}\right)$. For both curves: $a=2, b=3, d=1, \delta=0$.
and the travelling solutions of Equation (19) are obtained as
$U_{ \pm}\left(\xi_{n}\right)=-\frac{b}{2} \pm \frac{\sinh (d)}{2} \sqrt{b^{2}-4 a} \operatorname{csch}\left(\xi_{n}\right)$,
Where $\xi_{n}=n d+\frac{\sinh (d)}{2}\left(4 a-b^{2}\right) t+\delta$. The portraits of solutions (20) for $U_{ \pm}\left(\xi_{n}\right)$ are displayed in Figure 2(a) and (b).

## Third set

$c=\left(4 a-b^{2}\right) \tanh \left(\frac{d}{2}\right)$,
$\alpha=-\frac{b}{2}$,
$A= \pm \frac{1}{2} \sqrt{b^{2}-4 a} \tanh \left(\frac{d}{2}\right)$,
$B= \pm \frac{1}{2} \sqrt{b^{2}-4 a} \tanh \left(\frac{d}{2}\right)$,
Finally, third set admits the following two types:

$$
\begin{equation*}
U_{1 \pm}\left(\xi_{n}\right)=-\frac{b}{2}-\frac{1}{2} \tanh \left(\frac{d}{2}\right) \sqrt{b^{2}-4 a}\left(\operatorname{coth}\left(\xi_{n}\right) \pm \operatorname{csch}\left(\xi_{n}\right)\right), \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
U_{2 \pm}\left(\xi_{n}\right)=-\frac{b}{2}+\frac{1}{2} \tanh \left(\frac{d}{2}\right) \sqrt{b^{2}-4 a}\left(\operatorname{coth}\left(\xi_{n}\right) \pm \operatorname{csch}\left(\xi_{n}\right)\right), \tag{23}
\end{equation*}
$$

Where $\xi_{n}=n d+\left(4 a-b^{2}\right) \tanh \left(\frac{d}{2}\right) t+\delta$. The behaviors of solutions (22) and (23) for $U_{1+}\left(\xi_{n}\right)$ and $U_{2+}\left(\xi_{n}\right)$ are shown in Figure 3(a) and (b) respectively. The solutions given for the second and the third sets appear to be new.

## CONCLUSION

The basic goal of this work, is to provide a new trial travelling solution to build the exact solutions to the nonlinear lattice equation. Two types of functions are used to find the exact solutions, which are named the coth-function and the coth-csch function methods. Eight variants of travelling wave solutions are obtained. The present method provides a reliable technique that requires less work if compared with the difficulties arising from computational aspect. The main advantage of this method is the flexibility to give exact solutions to nonlinear problems without linearization. We may conclude that, this method can also be extended to other high-dimensional nonlinear phenomena. It will be then interesting to study more general systems. These points will be investigated in a future research.

## Conflict of Interest

The authors have not declared any conflict of interest.

## ACKNOWLEDGEMENTS

Authors want to thank Pr W. Higg and Dr J. Karim for their valuable comments. The project is supported by Ministère de l'Enseignement et de la Recherche Scientifique (M.E.R.S): PNR n ${ }^{\circ}$ 30/15/2011.


Figure 3. The graphs show the wave solution of $U_{1 \pm}\left(\xi_{n}\right)$ and $U_{2 \pm}\left(\xi_{n}\right)$ in (22) and (23) respectively: (a) solution $u_{1+}(n, t)=U_{1+}\left(\xi_{n}\right),(\mathrm{b})$ solution $u_{2+}(n, t)=U_{2+}\left(\xi_{n}\right)$. For both curves: $a=2, b=3, d=1, \delta=0$.

## REFERENCES

He JH, Wu XH (2006a). Exp-function method for nonlinear wave equations. Chaos Solitons Fractals. 30(3):700-708. http://dx.doi.org/10.1016/j.chaos.2006.03.020
He JH, Wu XH (2006b). Construction of solitary solution and compacton-like solution by variational iteration method. Chaos Solitons Fractals. 29(1):108-113. http://dx.doi.org/10.1016/j.chaos.2005.10.100
Ibrahim RS, El-Kalaawy OH (2007). Extended tanh-function method and reduction of nonlinear Schrödinger-type equations to a quadrature. Chaos Solitons Fractals. 31(4):1001-1008. http://dx.doi.org/10.1016/j.chaos.2005.10.055
Jia L, Liu Q, Ma Z (2014). A good approximation of modulated amplitude waves in Bose-Einstein condensates. Communications in Nonlinear Sci. Numer. Simul. 19(8):2715-2723. http://dx.doi.org/10.1016/j.cnsns.2013.12.034
Khelii N, Bensalah N, Saidi H, Zerarka A (2006). Artificial perturbation for solving the Korteweg-de Vries equation. J. Zhejiang Univ. SCI. A 7(12):2079-2082.
Korteweg DJ, de Vries G (1995). On the change of the form of long waves advancing in a rectangular canal and on a new type of long stationary waves. Philos. Mag. P. 39.
Liu Q, Qian D (2011). Modulated amplitude waves with nonzero phases in Bose-Einstein condensates. J. Math. Phys. 52(8):082702. http://dx.doi.org/10.1063/1.3623415
Lü X (2014a). New bilinear Backlund transformation with multisoliton solutions for the $(2+1)$-dimensional Sawada-Kotera model. Nonlinear Dyn. 76:161. http://dx.doi.org/10.1007/s11071-013-1118-y
Lü X (2014b), Bright-soliton collisions with shape change by intensity redistribution for the coupled Sasa-Satsuma system in the optical fiber communications, Communications in Nonlinear Science and Numerical Simulation, 19:3969. http://dx.doi.org/10.1016/j.cnsns.2014.03.013
Lü X, Peng M (2013a). Nonautonomous motion study on accelerated and decelerated solitons for the variable-coefficient Lenells-Fokas model. Chaos. 23:013122. http://dx.doi.org/10.1063/1.4790827
Lü X, Peng M (2013b). Systematic construction of infinitely many conservation laws for certain nonlinear evolution equations in mathematical physics. Commun. Nonlinear Sci. Numer. Simulat. 18:2304. http://dx.doi.org/10.1016/j.cnsns.2012.11.006
Lü X, Peng M (2013c). Painleve-integrability and explicit solutions of the general two-coupled nonlinear Schrodinger system in the optical fiber $\begin{array}{ll}\text { communications. Nonlinear } & \text { Dyn. 73:405. }\end{array}$ http://dx.doi.org/10.1007/s11071-013-0795-x
Lü X (2013). Soliton behavior for a generalized mixed nonlinear Schrodinger model with N-fold Darboux transformation. Chaos. 23:033137. http://dx.doi.org/10.1063/1.4821132

Lü X, Tian B, Zhang HQ, Xu T, Li H (2010). Integrability study on the generalized ( $2+1$ )-dimensional variable-coefficient Gardner model with symbolic computation. Chaos. 20:043125. http://dx.doi.org/10.1063/1.3494154
Malfliet W (1996a). The tanh method: I. Exact solutions of nonlinear evolution and wave equations. Phys. Scr. 54:563-568. http://dx.doi.org/10.1088/0031-8949/54/6/003
Malfliet W (1996b). The tanh method:II.Perturbation technique for conservative systems. Phys. Scr. 54:569-575. http://dx.doi.org/10.1088/0031-8949/54/6/004
Odibat ZM, Momani S (2006). Application of variational iteration method to nonlinear differential equations of fractional order. Int. J. Nonlinear Sci. Numer. Simul. 7(1):27-34. http://dx.doi.org/10.1515/IJNSNS.2006.7.1.27
Özis T, Yıldırım A (2008). Reliable analysis for obtaining exact soliton solutions of nonlinear Schrödinger (NLS) equation. Chaos Solitons Fractals. 38:209-212. http://dx.doi.org/10.1016/j.chaos.2006.11.006
Rafei M, Ganji DD (2006). Explicit solutions of Helmholtz equation and fifth-order KdV equation using homotopy perturbation method. Int. J. Nonlinear Sci. Numer. Simul. 7(3):321-328. http://dx.doi.org/10.1515/JJNSNS.2006.7.3.321
Wang DS, Li H (2008). Symbolic computation and non-travelling wave solutions of $(2+1)$-dimensional nonlinear evolution equations. Chaos Solitons Fractals.

38:383-390. http://dx.doi.org/10.1016/j.chaos.2007.07.062
Wazwaz AM (2004). The tanh method for travelling wave solutions of nonlinear equations. Appl. Math. Comput. 154(3):713-23. http://dx.doi.org/10.1016/S0096-3003(03)00745-8
Xu LP, Zhang JL (2007). Exact solutions to two higher order nonlinear Schrödinger equations. Chaos Solitons Fractals. 31:937-942. http://dx.doi.org/10.1016/j.chaos.2005.10.063
Yu YX (2007). Rational formal solutions of hybrid lattice equation. Appl. Math. Comput. 186:474-482. http://dx.doi.org/10.1016/j.amc.2006.07.112
Zerarka A, Foester VG (2005). Separation method for solving the generalized Korteweg --de Vries equation. Comm. Nonlinear. Sci. Numer. Simul. 10:217-225. http://dx.doi.org/10.1016/j.cnsns.2003.09.003
Zhu SD (2007a). Exp-function method for the hybrid-lattice systems. Int. J. Nonlinear Sci. Numer. Simul. 8:461-464. http://dx.doi.org/10.1515/IJNSNS.2007.8.3.465
Zhu SD (2007b). Exp-function method for the discrete mkdv lattice. Int. J. Nonlinear Sci. Numer. Simul. 8:465-468. http://dx.doi.org/10.1515/IJNSNS.2007.8.3.461
Zhu SD (2008). Discrete (2+1)-dimensional toda lattice equation via exp-function method. Phys. Lett. A 372(5):654-657. http://dx.doi.org/10.1016/j.physleta.2007.07.085


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