## Full Length Research Paper

# Spacelike ${ }_{B_{2}}$-slant helix in Minkowski 4-space $E_{E_{1}^{4}}$ <br> Mehmet Önder*, Hüseyin Kocayiğit and Mustafa Kazaz <br> Department of Mathematics, Faculty of Science and Art, Celal Bayar University, 45047 Manisa, Turkey. 

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#### Abstract

In this paper, we give the characterizations of spacelike $B_{2}$-slant helix by means of curvatures of the spacelike curve in Minkowski 4 - space. Furthermore, we give the integral characterization of the spacelike $B_{2}$-slant helix.


Key words: Minkowski 4 - space, spacelike $B_{2}$-slant helix, Frenet frame.

## INTRODUCTION

Helix is one of the most fascinating curves in Science and Nature. Helices can be seen in many subjects of Science such as nanosprings, carbon nanotubes, $\alpha$-helices, DNA double and collagen triple helix (a DNA molecule has a form of well - known right-handed double helix, wherein two heteropolymer chains are wound around each other. The double helical structure is believed to be the structure of minimum free energy under the normal physiological conditions.), lipid bilayers, bacterial flagella in Escherichia coli and Salmonella, aerial hyphae in actinomycetes, bacterial shape in spirochetes, horns, tendrils, vines, screws, springs, helical staircases and sea shells (helico-spiral structures) (Chouaieb et al., 2006; Lucas Amand and Lambin, 2005; Watson and Crick, 1953). Furthermore, in the fields of computer aided design and computer graphics, helices can be used for the tool path description, the simulation of kinematic motion or the design of highways, etc. (Yang, 2003).
A curve of constant slope or general helix in Euclidean 3 -space $E^{3}$ is defined by the property that the tangent makes a constant angle with a fixed straight line which is called the axis of the general helix (Barros, 1997). A classical result stated by M. A. Lancret in 1802 and first proved by B. de Saint Venant in 1845 is: A necessary and sufficient condition that a curve be a general helix is that the ratio of the first curvature to the second curvature be constant that is $k_{1} / k_{2}$ is constant along the curve, where $k_{1}$ and $k_{2}$ denote the first and second curvatures

[^0]of the curve, respectively (Scofield, 1995). Analogue to that Magden (1993) has given a characterization for a curve $x(s)$ to be a helix in Euclidean 4-space $E^{4}$. He characterizes a helix if the function
$\frac{k_{1}^{2}}{k_{2}^{2}}+\left[\frac{1}{k_{3}} \frac{d}{d s}\left(\frac{k_{1}}{k_{2}}\right)\right]^{2}$
is constant where $k_{1}, k_{2}$ and $k_{3}$ are first, second and third curvatures of Euclidean curve $x(s)$, respectively, and they are nowhere zero. Corresponding characterizations of time like helices in Minkowski 4space $E_{1}^{4}$ were given by Kocayigit and Onder (2007). Latterly Camci et al., (2009) have given some characterizations for a non - degenerate curve to be a generalized helix by using its harmonic curvatures.
Recently, Izumiya and Takeuchi (2004) have introduced the concept of slant helix by saying that the normal lines of the curve make a constant angle with a fixed direction and they have given a characterization of slant helix in Euclidean 3 -space $E^{3}$ by the fact that the function
$\frac{k_{1}^{2}}{\left(k_{1}^{2}+k_{2}^{2}\right)^{3 / 2}}\left(\frac{k_{2}}{k_{1}}\right)^{\prime}$
is constant. After them, Kula and Yayli (2005)
investigated spherical images, the tangent indicatrix and the binormal indicatrix of a slant helix and they obtained that the spherical images are spherical helices. Analogue to the definition of slant helix, Onder et al. (2008) have defined $B_{2}$-slant helix in Euclidean 4-space $E^{4}$ by saying that the second binormal vector of the curve make a constant angle with a fixed direction and they have given some characterizations of $B_{2}$-slant helix in Euclidean 4-space $E^{4}$.
In this paper, we consider spacelike $B_{2}$ - slant helix in Minkowski 4-space $E_{1}^{4}$ and we give some characterizations and also the integral characterization of spacelike $B_{2}$-slant helix.

## MATERIALS AND METHODS

Minkowski space-time $E_{1}^{4}$ is a Euclidean space $E^{4}$ provided with the standard flat metric given by

$$
\langle,\rangle=-d x_{1}^{2}+d x_{2}^{2}+d x_{3}^{2}+d x_{4}^{2}
$$

where $\left(x_{1}, x_{2}, x_{3}, x_{4}\right)$ is a rectangular coordinate system in $E_{1}^{4}$.
Since $\langle$,$\rangle is an indefinite metric, recall that a vector v \in E_{1}^{4}$ can have one of three causal characters; it can be spacelike if $\langle v, v\rangle>0$ or $v=0$, time like if $\langle v, v\rangle<0$ and null (light like) if $\langle v, v\rangle=0$ and $v \neq 0$. Similarly, an arbitrary curve $x(s)$ in $E_{1}^{4}$ can locally be spacelike, time like or null (light like), if all of its velocity vectors $x^{\prime}(s)$ are respectively spacelike, time like or null (light like). Also recall that the pseudo-norm of an arbitrary vector $v \in E_{1}^{4}$ is given by $\|v\|=\sqrt{|\langle v, v\rangle|}$. Therefore $v$ is a unit vector if $\langle v, v\rangle= \pm 1$. The velocity of the curve $x(s)$ is given by $\left\|x^{\prime}(s)\right\|$. Next, vectors $v, w$ in $E_{1}^{4}$ are said to be orthogonal if $\langle v, w\rangle=0$. We say that a time like vector is future pointing or past pointing if the first compound of the vector is positive or negative, respectively. Denote by $\left\{T(s), N(s), B_{1}(s), B_{2}(s)\right\}$ the moving Frenet frame along the curve $x(s)$ in the space $E_{1}^{4}$. Then $T, N, B_{1}, B_{2}$ are the tangent, the principal normal, the first binormal and the second binormal fields, respectively. A timelike (resp. spacelike) curve $x(s)$ is said to be parameterized by a pseudo - arc length parameter $s$, that is $\left\langle x^{\prime}(s), x^{\prime}(s)\right\rangle=-1$ (resp. $\left\langle x^{\prime}(s), x^{\prime}(s)\right\rangle=1$ ). Let $x(s)$ be a spacelike curve in Minkowski space-time $E_{1}^{4}$, parameterized by arc length function of $S$. Then for the curve $x(s)$ the following Frenet equation is given as follows:

$$
\left[\begin{array}{l}
T^{\prime}  \tag{1}\\
N^{\prime} \\
B_{1}^{\prime} \\
B_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cccc}
0 & k_{1} & 0 & 0 \\
-\varepsilon_{1} k_{1} & 0 & k_{2} & 0 \\
0 & \varepsilon_{2} k_{2} & 0 & k_{3} \\
0 & 0 & \varepsilon_{1} k_{3} & 0
\end{array}\right]\left[\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right]
$$

where

$$
\langle T, T\rangle=1,\langle N, N\rangle=\varepsilon_{1},\left\langle B_{2}, B_{2}\right\rangle=\varepsilon_{2}
$$

$$
\left\langle B_{1}, B_{1}\right\rangle=-\varepsilon_{1} \varepsilon_{2} \varepsilon_{1}= \pm 1, \varepsilon_{2}= \pm 1
$$

and recall that the functions $k_{1}=k_{1}(s), k_{2}=k_{2}(s)$ and $k_{3}=k_{3}(s)$ are called the first, the second and the third curvature of the spacelike curve $x(s)$, respectively and we will assume throughout this work that all the three curvatures satisfy $k_{i}(s) \neq 0,1 \leq i \leq 3$. Here the signs of $\varepsilon_{1}$ and $\varepsilon_{2}$ are changed by a rule. The signature rule between $\varepsilon_{1}$ and $\varepsilon_{2}$ can be given as follows

or

| if | $\varepsilon_{2}+1$ | then | $\varepsilon_{1}+1$ or -1 |
| :--- | :--- | :--- | :--- |
|  | -1 |  | +1 |

For the obvious forms of the Frenet equations in (1) we refer to the reader to see Walrave (1995).

## RESULTS AND DISCUSSION

In this section, we give the definition and the characterizations of spacelike $B_{2}$-slant helix.
Let $x: I \subset I R \rightarrow E_{1}^{4}$ be a unit speed spacelike curve with nonzero curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ and let $\left\{T, N, B_{1}, B_{2}\right\}$ denotes the Frenet frame of the curve $x(s)$. We call $x(s)$ as spacelike $B_{2}$-slant helix if its second binormal unit vector $B_{2}$ makes a constant angle with a fixed direction in a unit vector $U$; that is
$\left\langle B_{2}, U\right\rangle=$ constant
along the curve. By differentiation (2) with respect to $s$
and using the Frenet formulae (1) we have $\left\langle\varepsilon_{1} k_{3} B_{1}, U\right\rangle=0$.

Therefore $U$ is in the subspace $\operatorname{Sp}\left\{T, N, B_{2}\right\}$ and can be written as follows

$$
\begin{equation*}
U=a_{1}(s) T(s)+a_{2}(s) N(s)+a_{3}(s) B_{2}(s) \tag{3}
\end{equation*}
$$

where
$a_{1}=\langle U, T\rangle, \varepsilon_{1} a_{2}=\langle U, N\rangle$,
$\varepsilon_{2} a_{3}=\left\langle U, B_{2}\right\rangle=$ constant
Since $U$ is unit, we have
$a_{1}^{2}+\varepsilon_{1} a_{2}^{2}+\varepsilon_{2} a_{3}^{2}=M$.
Here $M$ is $+1,-1$ or 0 depending if $U$ is spacelike, timelike or lightlike, respectively. The differentiation of (3) gives

$$
\begin{aligned}
& \left(\frac{d a_{1}}{d s}-\varepsilon_{1} a_{2} k_{1}\right) T+\left(\frac{d a_{2}}{d s}-a_{1} k_{1}\right) N \\
& +\left(a_{2} k_{2}+\varepsilon_{1} a_{3} k_{3}\right) B_{1}+a_{3}^{\prime} B_{2}=0
\end{aligned}
$$

and from this equation we get
$\left.\begin{array}{l}a_{2}=-\varepsilon_{1} \frac{k_{3}}{k_{2}} a_{3}=\varepsilon_{1} \frac{1}{k_{1}} \frac{d a_{1}}{d s}, \\ \frac{d a_{2}}{d s}=-a_{1} k_{1}, a_{3}^{\prime}=0\end{array}\right\}$
Since $\frac{d a_{2}}{d s}=-a_{1} k_{1}$ and
$\frac{d a_{2}}{d s}=-\varepsilon_{1} \frac{k_{1}^{\prime}}{k_{1}^{2}} \frac{d a_{1}}{d s}+\frac{\varepsilon_{1}}{k_{1}} \frac{d^{2} a_{1}}{d s^{2}}$
we find the second order linear differential equation in $a_{1}$ given by
$\varepsilon_{1} \frac{d^{2} a_{1}}{d s^{2}}-\varepsilon_{1} \frac{k_{1}^{\prime}}{k_{1}} \frac{d a_{1}}{d s}+a_{1} k_{1}^{2}=0$.
If we change variables in the above equation as
$t=\int_{0}^{s} k_{1}(s) d s$ then we get
$\frac{d^{2} a_{1}}{d t^{2}}+\varepsilon_{1} a_{1}=0$.
This equation has two solutions: If $\varepsilon_{1}=+1$ then the solution is
$a_{1}=A \cos \int_{0}^{s} k_{1}(s) d s+B \sin \int_{0}^{s} k_{1}(s) d s$
and if $\varepsilon_{1}=-1$ then the solution is
$a_{1}=A \cosh \int_{0}^{s} k_{1}(s) d s+B \sinh \int_{0}^{s} k_{1}(s) d s$
where $A$ and $B$ are constant.
Assume that $\varepsilon_{1}=-1$ and consider the solution (8). From (5) and (8) we have
$a_{2}=\frac{k_{3}}{k_{2}} a_{3}=-A \sinh \int_{0}^{s} k_{1}(s) d s-B \cosh \int_{0}^{s} k_{1}(s) d s$

$$
\begin{aligned}
a_{1}=-\frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime} a_{3}= & A \cosh \int_{0}^{s} k_{1}(s) d s \\
& +B \sinh \int_{0}^{s} k_{1}(s) d s
\end{aligned}
$$

From these equations it follows that

$$
\begin{align*}
& A=a_{3}\left(\frac{k_{3}}{k_{2}} \sinh \int_{0}^{s} k_{1}(s) d s-\frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime} \cosh \int_{0}^{s} k_{1}(s) d s\right)  \tag{9}\\
& B=a_{3}\left(-\frac{k_{3}}{k_{2}} \cosh \int_{0}^{s} k_{1}(s) d s+\frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime} \sinh \int_{0}^{s} k_{1}(s) d s\right) \tag{10}
\end{align*}
$$

Hence, using (9) and (10) we get
$B^{2}-A^{2}=\left[\left(\frac{k_{3}}{k_{2}}\right)^{2}-\frac{1}{k_{1}^{2}}\left(\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right)^{2}\right] a_{3}^{2}=$ constant

So that
$\left(\frac{k_{3}}{k_{2}}\right)^{2}-\frac{1}{k_{1}^{2}}\left(\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right)^{2}=$ constant $:=m$.
From (4), (9), (10) and (11) we have
$B^{2}-A^{2}=a_{3}^{2} m=M-1$.
Thus, the sign of the constant $m$ agrees with the sign of $B^{2}-A^{2}$. So, if $U$ is timelike or light like then $m$ is negative. If $U$ is spacelike then $m=0$. Then we can give the following corollary.

Result 1: Let $x(s)$ be a spacelike $B_{2}$-slant helix with timelike principal normal $N$ in Minkowski 4-space $E_{1}^{4}$ and $U$ be a unit constant vector which makes a constant angle with the second binormal $B_{2}$. Then the vector $U$ is spacelike if and only if there exist a constant $K$ such that
$\frac{k_{3}}{k_{2}}(s)=K \exp \left(\int_{0}^{s} k_{1}(t) d t\right)$.
When $\varepsilon_{1}=+1$, by using (7) with similar calculations as above, we get that the spacelike curve $x(s)$ is a spacelike $B_{2}$-slant helix if and only if
$\left(\frac{k_{3}}{k_{2}}\right)^{2}+\frac{1}{k_{1}^{2}}\left(\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right)^{2}=$ constant .
Thus, using (11) and (12), we can characterize the spacelike $B_{2}$-slant helix $x(s)$ by the fact that
$\left(\frac{k_{3}}{k_{2}}\right)^{2}+\varepsilon_{1} \frac{1}{k_{1}^{2}}\left(\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right)^{2}=$ constant.
Conversely, if the condition (13) is satisfied for a regular spacelike curve we can always find a constant vector $U$ which makes a constant angle with the second binormal $B_{2}$ of the curve.
Consider the unit vector $U$ defined by
$U=\left[\varepsilon_{1} \frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime} T-\varepsilon_{1} \frac{k_{3}}{k_{2}} N+B_{2}\right]$.
By taking account of the differentiation of (13), differentiation of $U$ gives that $\frac{d U}{d s}=0$, this means that $U$ is a constant vector. So that, we can give the following theorem:

Theorem 1: A unit speed spacelike curve $x: I \subset I R \rightarrow E_{1}^{4}$ with nonzero curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ is a spacelike $B_{2}$-slant helix if and only if the following condition is satisfied,
$\left(\frac{k_{3}}{k_{2}}\right)^{2}+\varepsilon_{1} \frac{1}{k_{1}^{2}}\left(\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right)^{2}=$ constant .
From Theorem 1 one can easily see that the constant function in Theorem 1 is independent of $\varepsilon_{2}$. So, we can give the following corollary.

Result 2: The characterizations of the spacelike $B_{2}$-slant helix are independent of the Lorentzian causal character of the second binormal vector $B_{2}$. It is only related to the Lorentzian causal character of the unit principal normal vector $N$.

Now, we give another characterization of spacelike $B_{2}$ slant helix in Minkowski 4-space.
Let assume that $x: I \subset I R \rightarrow E_{1}^{4}$ is a spacelike $B_{2}$ slant helix. Then, Theorem 1 is satisfied. By differentiating (13) with respect to $s$ we get

$$
\begin{equation*}
\left(\frac{k_{3}}{k_{2}}\right) \frac{d}{d s}\left(\frac{k_{3}}{k_{2}}\right)+\frac{\varepsilon_{1}}{k_{1}} \frac{d}{d s}\left(\frac{k_{3}}{k_{2}}\right) \frac{d}{d s}\left[\frac{1}{k_{1}} \frac{d}{d s}\left(\frac{k_{3}}{k_{2}}\right)\right]=0 \tag{14}
\end{equation*}
$$

and hence

$$
\begin{equation*}
\frac{\varepsilon_{1}}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime}=-\frac{\left(\frac{k_{3}}{k_{2}}\right)\left(\frac{k_{3}}{k_{2}}\right)^{\prime}}{\left[\frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right]^{\prime}} . \tag{15}
\end{equation*}
$$

If we write

$$
\begin{equation*}
f(s)=-\frac{\left(\frac{k_{3}}{k_{2}}\right)\left(\frac{k_{3}}{k_{2}}\right)^{\prime}}{\left[\frac{1}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right]^{\prime}} . \tag{16}
\end{equation*}
$$

Then

$$
\begin{equation*}
f(s) k_{1}=\varepsilon_{1}\left(\frac{k_{3}}{k_{2}}\right)^{\prime} \tag{17}
\end{equation*}
$$

From (14) it can be written

$$
\begin{equation*}
\left[\frac{\varepsilon_{1}}{k_{1}}\left(\frac{k_{3}}{k_{2}}\right)^{\prime}\right]^{\prime}=-k_{1} \frac{k_{3}}{k_{2}} . \tag{18}
\end{equation*}
$$

By using (17) and (18) we have

$$
\begin{equation*}
\frac{d}{d s} f(s)=-k_{1} \frac{k_{3}}{k_{2}} \tag{19}
\end{equation*}
$$

Conversely,

$$
\text { let } \quad f(s) k_{1}=\varepsilon_{1}\left(\frac{k_{3}}{k_{2}}\right)^{\prime}
$$

and $\frac{d}{d s} f(s)=-k_{1} \frac{k_{3}}{k_{2}}$. If we define a unit vector $U$ by
$U=-f(s) T+\varepsilon_{1} \frac{k_{3}}{k_{2}} N-B_{2}$

We have that $U$ and $\left\langle B_{2}, U\right\rangle$ are constants. So, we have the following theorem:

Theorem 2: A unit speed spacelike curve $x: I \subset I R \rightarrow E_{1}^{4}$ with nonzero curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ is a $B_{2}$-slant helix if and only if there exists a $C^{2}$-function $f$ such that
$f k_{1}=\varepsilon_{1} \frac{d}{d s}\left(\frac{k_{3}}{k_{2}}\right), \frac{d}{d s} f(s)=-k_{1} \frac{k_{3}}{k_{2}}$

Now, we give the integral characterization of the spacelike $B_{2}$-slant helix.
Suppose that, the unit speed spacelike curve $x: I \subset I R \rightarrow E_{1}^{4}$ with nonzero curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ is a spacelike $B_{2}$-slant helix. Then Theorem 2 is satisfied. Let us define $C^{2}$-function $\varphi$ and $C^{1}$ functions $m(s)$ and $n(s)$ by
$\varphi=\varphi(s)=\int_{0}^{s} k_{1}(s) d s$
$\left.\begin{array}{l}m(s)=\frac{k_{3}}{k_{2}} \eta(\varphi)+f(s) \mu(\varphi) \\ n(s)=\frac{k_{3}}{k_{2}} \mu(\varphi)-\varepsilon_{1} f(s) \eta(\varphi)\end{array}\right\}$
Where $\quad \eta(\varphi)=\cosh \varphi, \quad \mu(\varphi)=\sinh \varphi \quad$ if $\varepsilon_{1}=-1 ;$ and $\eta(\varphi)=\cos \varphi, \quad \mu(\varphi)=\sin \varphi \quad$ if $\varepsilon_{1}=+1$. If we differentiate equations (23) with respect to $s$ and take account of (22) and (21) we find that $m^{\prime}=0$ and $n^{\prime}=0$. Therefore, $m(s)=C, n(s)=D$ are constants. Now substituting these in (23) and solving the resulting equations for $\frac{k_{3}}{k_{2}}$, we get
$\frac{k_{3}}{k_{2}}=C \eta(\varphi)+D \mu(\varphi)$.
Conversely if (24) holds then from the equations in (23) we get

$$
f=\varepsilon_{1}(C \mu(\varphi)-D \eta(\varphi))
$$

which satisfies the conditions of Theorem 2? So, we have the following theorem:

Theorem 3: A unit speed spacelike curve $x: I \subset I R \rightarrow E_{1}^{4}$ with nonzero curvatures $k_{1}(s), k_{2}(s)$ and $k_{3}(s)$ is a spacelike $B_{2}$-slant helix if and only if the following condition is satisfied

$$
\frac{k_{3}}{k_{2}}=C \eta(\varphi)+D \mu(\varphi)
$$

where $C$ and $D$ are constants, $\eta(\varphi)=\cosh \varphi$,
$\mu(\varphi)=\sinh \varphi \quad$ if $\quad \varepsilon_{1}=-1 ; \quad$ and $\quad \eta(\varphi)=\cos \varphi$, $\mu(\varphi)=\sin \varphi$ if $\varepsilon_{1}=+1$.

## Conclusions

In this paper, the spacelike $B_{2}$-slant helix is defined and the characterizations of the spacelike $B_{2}$-slant helix are given in Minkowski 4- space $E_{1}^{4}$. It is shown that a spacelike curve $x: I \subset I R \rightarrow E_{1}^{4}$ is a $B_{2}$-slant helix if an equation holds between the first, second and third curvatures of the curve. Furthermore, the integral characterization of the spacelike $B_{2}$-slant helix is given.

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