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Scientific Research and Essays

Full Length Research Paper

### The improved generalized Riccati equation mapping method and its application for solving a nonlinear partial differential equation (PDE) describing the dynamics of ionic currents along microtubules

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In this paper we apply the improved Riccati equation mapping method to construct many families of exact solutions of a nonlinear partial differential equation involving parameters of a special interest in nanobiosciences and biophysics which describe a model of microtubules as nonlinear RLC transmission lines. As results, we can successfully recover the previously known results that have been found using other methods. This method is straightforward and concise, and it can be applied to other nonlinear PDEs in mathematical physics. Comparison between our new results and the well-known results are given. Some comments on the well-known results are also presented at the end of this article.

**Key words:** Improved Riccati equation mapping method, exact traveling wave solutions, nonlinear partial differential equations (PDEs) of microtubules, Nonlinear RLC transmission lines.

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### INTRODUCTION

In recent years, the exact traveling wave solutions for nonlinear partial differential equations (PDEs) has been investigated by many authors who are interested in non linear physical phenomena. Many powerful methods have been presented, such as the inverse scattering transform method (Ablowitz and Clarkson, 1991), the Hirota's bilinear method (Hirota, 1971), the Painleve expansion method (Weiss et al., 1983; Kudryashov, 1988, 1990, 1991), the Backlund truncated method (Miura, 1978; Rogers and Shadwick, 1982), the exp-function method (He and Wu, 2006; Yusufoglu, 2008; Zhang, 2008; Bekir, 2009, 2010; Aslan, 2011), the tanh-function method (Abdou, 2007; Fan, 2000; Zhang and Xia, 2008; Yusufoglu and Bekir, 2008), the Jacobi elliptic function

\*Corresponding author. E-mail: e.m.e.zayed@hotmail.com Author(s) agree that this article remain permanently open access under the terms of the <u>Creative Commons Attribution</u> <u>License 4.0 International License</u> method (Chen and Wang, 2005; Liu et al., 2001; Lu, 2005), the (G'/G)-expansion method (Wang et al., 2008; Zhang et al., 2008; Zayed, 2009, 2010; Bekir, 2008; Ayhan and Bekir, 2012, Kudryashov, 2010a,b; Aslan, 2010, 2011, 2012a,b), the generalized Riccati equation mapping method (Zhu 2008; Zayed and Arnous, 2013, Zayed et al., 2013), and so on.

In the present paper, we shall use the improved Riccati equation mapping method to find the exact solutions of a nonlinear PDE of nanobiosciences. The main idea of this method is that the traveling wave solutions of nonlinear equations can be expressed by polynomials in Q, where  $Q = Q(\xi)$  satisfies the generalized Riccati equation  $Q' = r + pQ + qQ^2$  where  $\xi = kx + \omega t$ , where  $r, p, k, \omega$  and q are constants. The degree of this polynomial can be determined by considering the homogeneous balance between the highest order derivatives and the nonlinear terms appearing in the given nonlinear equation, the coefficients of this polynomial can obtained by solving a set of algebraic equations resulted from the process of using the proposed method.

The objective of this paper is to apply the improved Riccati equation mapping method for finding many families of exact traveling wave solutions of the following nonlinear PDE of special interest in nanobiosciences, namely, the transmission line models for microtubules as nonlinear RLC transmission line (Sekulic et al., 2011a, Sataric et al., 2010):

$$R_2 C_0 L^2 u_{xxt} + L^2 u_{xx} + 2R_1 C_0 \delta u u_t - R_1 C_0 u_t = 0$$
(1)

where  $R_1 = 10^9 \Omega$  and  $R_2 = 7 \times 10^6 \Omega$  stand for longitudinal and transversal component of resistance of an Elementary rings and parameter  $\delta(\delta < 1)$  describes nonlinearity of ER capacitor in MT. Here  $L = 8 \times 10^{-9} m$ while  $C_0 = 1.8 \times 10^{-15} F$  is the total maximal capacitance of the ER. The physical details of the derivation of Equation (1) can be elaborated in Sataric et al. (2010). For further references about electrical models of microtubules, see for example llic et al. (2009), Sekulic et al. (2011b, 2012), Sataric et al. (2009); Freedman et al. (2010), and Sekulic and Sataric (2012). Recently, Equation (1) has been discussed in (Sekulic et al. 2011a) by using the modified extended tanh-function method, where its exact solutions have been found.

The rest of this paper can be organized as follows: First is description of the improved generalized Riccati equation method. Many families of exact traveling wave solutions for Equation (1) are next obtained. This is followed by illustrations on physical explanations for some obtained results. Thereafter, conclusions and comments on Sekulic et al. (2011a) as well as comparison between our new results and the well-known results obtained in Sekulic et al. (2011a) are investigated.

# Description of the improved generalized Riccati equation mapping method

We suppose that a nonlinear PDE is in the following from:

$$P(u, u_x, u_t, u_{xx}, u_{tt}, ...) = o$$
<sup>(2)</sup>

where u = u(x,t) is an unknown function , *P* is a polynomial in u = u(x,t) and its partial derivatives in which the highest order derivatives and nonlinear terms are involved. Let us now give the main steps for solving Equation (2.1) using the improved Riccati equation mapping method (Zhu 2008; Zayed and Arnous, 2013; Zayed et al, 2013):

Step 1: We look for its traveling wave solution in the form

$$u(x,t) = u(\xi), \quad \xi = kx + \omega t \tag{3}$$

where  $k, \omega$  are constants. Substituting (3) into Equation (2) gives the nonlinear ODE for  $u(\xi)$  as follows:

$$H(u,u',u'',...) = 0$$
(4)

where H is a polynomial in  $u(\xi)$  and its total derivatives

$$u', u'', u''', \dots$$
 such that  $u' = \frac{du}{d\xi}, u'' = \frac{d^2u}{d\xi^2}, \dots$ 

**Step 2:** We suppose that the solution of the ODE (4) can be expressed as follows:

$$u(\xi) = \sum_{i=-m}^{m} a_i Q^i(\xi) , \qquad (5)$$

where  $a_i (i = 0, \pm 1, \pm 2, ..., \pm m)$  are constants to be determined later such as  $a_m \neq 0$  or  $a_{-m} \neq 0$  and  $Q = Q(\xi)$  is the solution of generalized Riccati equation

$$Q' = r + pQ + qQ^2 \tag{6}$$

where r, p and q are constants, such that  $q \neq 0$ .

**Step 3:** We determine the positive integer *m* in (5) by balancing the nonlinear terms and the highest order derivatives of  $u(\xi)$  in Equation (4).

**Step 4:** Substituting (5) and along with Equation (6) into Equation (4) and then equating all the coefficients of  $Q^i$  ( $i = 0, \pm 1, \pm 2, ..., \pm m$ ) to zero yield a system of

algebraic equations which can be solved by using the Maple or Mathematica to find the values of the constants  $a_i(-m,...,m)$  and  $k,\omega$ .

**Step 5:** It is well-known (Zhu 2008; Zayed and Arnous, 2013; Zayed et al., 2013) that Equation (6) has many families of solutions as follows:

**Type 1:** When  $\Delta = p^2 - 4qr > 0$  and  $pq \neq 0$  or  $qr \neq 0$  we have

$$\begin{split} \Phi_1(\xi) &= -\frac{1}{2q} [p + \sqrt{\Delta} \tanh(\frac{\sqrt{\Delta}}{2}\xi)], \\ \Phi_2(\xi) &= -\frac{1}{2q} [p + \sqrt{\Delta} \coth(\frac{\sqrt{\Delta}}{2}\xi)], \\ \Phi_3(\xi) &= -\frac{1}{2q} [p + \sqrt{\Delta} (\tanh(\sqrt{\Delta}\xi) \pm i \operatorname{sech}(\sqrt{\Delta}\xi))], \quad i\sqrt{-1} \\ \Phi_4(\xi) &= -\frac{1}{2q} [p + \sqrt{\Delta} (\coth(\sqrt{\Delta}\xi) \pm \operatorname{csch}(\sqrt{\Delta}\xi))], \\ \Phi_5(\xi) &= -\frac{1}{4q} [2p + \sqrt{\Delta} (\coth(\frac{\sqrt{\Delta}}{4}\xi) \pm \coth(\frac{\sqrt{\Delta}}{4}\xi))], \\ \Phi_6(\xi) &= \frac{1}{2q} [-p + \frac{\sqrt{\Delta}(A^2 + B^2)}{A \sinh(\sqrt{\Delta}\xi) + B}], \\ \Phi_7(\xi) &= \frac{1}{2q} [-p - \frac{\sqrt{\Delta}(B^2 - A^2)}{A \sinh(\sqrt{\Delta}\xi) + B}], \end{split}$$

where A and B are two non-zero real constants satisfying  $B^2$ - $A^2$ >0,

$$\begin{split} \Phi_{\mathfrak{s}}(\xi) &= \frac{2r\cosh(\frac{\sqrt{\Delta}}{2}\xi)}{\sqrt{\Delta}\sinh(\frac{\sqrt{\Delta}}{2}\xi) - p\cos h(\frac{\sqrt{\Delta}}{2}\xi)}, \\ \Phi_{\mathfrak{s}}(\xi) &= \frac{-2r\sinh(\frac{\sqrt{\Delta}}{2}\xi)}{p\sinh(\frac{\sqrt{\Delta}}{2}\xi) - \sqrt{\Delta}\cosh(\frac{\sqrt{\Delta}}{2}\xi)}, \\ \Phi_{10}(\xi) &= \frac{2r\cosh(\frac{\sqrt{\Delta}}{2}\xi)}{\sqrt{\Delta}\sinh(\sqrt{\Delta}\xi) - p\cosh(\sqrt{\Delta}\xi) \pm i\sqrt{\Delta}}, \quad i = \sqrt{-1} \\ \Phi_{11}(\xi) &= \frac{2r\sinh(\frac{\sqrt{\Delta}}{2}\xi)}{-p\sinh(\sqrt{\Delta}\xi) + \sqrt{\Delta}\cosh(\sqrt{\Delta}\xi) \pm \sqrt{\Delta}}, \\ \Phi_{12}(\xi) &= \frac{4r\sinh(\frac{\sqrt{\Delta}}{4}\xi)\cosh(\frac{\sqrt{\Delta}}{4}\xi)}{-2p\sinh(\frac{\sqrt{\Delta}}{4}\xi)\cosh(\frac{\sqrt{\Delta}}{4}\xi) + 2\sqrt{\Delta}\cosh^2(\frac{\sqrt{\Delta}}{2}\xi) - \sqrt{\Delta}}, \end{split}$$

**Type 2:** When  $\Delta = p^2 - 4qr < 0$  and  $pq \neq 0$  or  $qr \neq 0$  we have

$$\begin{split} \Phi_{13}(\xi) &= \frac{1}{2q} \left[ -p + \sqrt{-\Delta} \tan(\frac{\sqrt{-\Delta}}{2}\xi) \right], \\ \Phi_{14}(\xi) &= -\frac{1}{2q} \left[ p + \sqrt{-\Delta} \cot(\frac{\sqrt{-\Delta}}{2}\xi) \right], \\ \Phi_{15}(\xi) &= \frac{1}{2q} \left[ -p + \sqrt{-\Delta} (\tan(\sqrt{-\Delta}\xi) \pm \sec(\sqrt{-\Delta}\xi)) \right], \\ \Phi_{16}(\xi) &= -\frac{1}{2q} \left[ p + \sqrt{-\Delta} (\cot(\sqrt{-\Delta}\xi) \pm \csc(\sqrt{-\Delta}\xi)) \right], \\ \Phi_{17}(\xi) &= \frac{1}{4q} \left[ -2p + \sqrt{-\Delta} (\tan(\frac{\sqrt{-\Delta}}{4}\xi) - \cot(\frac{\sqrt{-\Delta}}{4}\xi)) \right], \\ \Phi_{18}(\xi) &= \frac{1}{2q} \left[ -p + \frac{\pm\sqrt{-\Delta}(A^2 - B^2) - A\sqrt{-\Delta} \cos(\sqrt{-\Delta}\xi)}{A\sin(\sqrt{-\Delta}\xi) + B} \right], \\ \Phi_{19}(\xi) &= \frac{1}{2q} \left[ -p - \frac{\pm\sqrt{-\Delta}(A^2 - B^2) - A\sqrt{-\Delta} \sin(\sqrt{-\Delta}\xi)}{A\sin(\sqrt{-\Delta}\xi) + B} \right], \end{split}$$

where A and B are two non-zero real constants satisfying  $A^2$ - $B^2$ >0,

$$\begin{split} \Phi_{20}(\xi) &= -\frac{2r\cos(\frac{\sqrt{-\Delta}}{2}\xi)}{\sqrt{-\Delta}\sin(\frac{\sqrt{-\Delta}}{2}\xi) + p\cos(\frac{\sqrt{-\Delta}}{2}\xi)},\\ \Phi_{21}(\xi) &= \frac{2r\sin(\frac{\sqrt{-\Delta}}{2}\xi)}{-p\sin(\frac{\sqrt{-\Delta}}{2}\xi) + \sqrt{-\Delta}\cos(\frac{\sqrt{-\Delta}}{2}\xi)},\\ \Phi_{22}(\xi) &= -\frac{2r\cos(\frac{\sqrt{-\Delta}}{2}\xi)}{\sqrt{-\Delta}\sin(\sqrt{-\Delta}\xi) + p\cos(\sqrt{-\Delta}\xi) \pm \sqrt{-\Delta}},\\ \Phi_{23}(\xi) &= \frac{2r\sin(\frac{\sqrt{-\Delta}}{2}\xi)}{-p\sin(\sqrt{-\Delta}\xi) + \sqrt{-\Delta}\cos(\sqrt{-\Delta}\xi) \pm \sqrt{-\Delta}},\\ \Phi_{24}(\xi) &= \frac{4r\sin(\frac{\sqrt{-\Delta}}{4}\xi)\cos(\frac{\sqrt{-\Delta}}{4}\xi)}{-2p\sin(\frac{\sqrt{-\Delta}}{4}\xi)\cos(\frac{\sqrt{-\Delta}}{4}\xi) + 2\sqrt{-\Delta}\cos^2(\frac{\sqrt{-\Delta}}{2}\xi) - \sqrt{-\Delta}} \end{split}$$

**Type 3:** When 
$$r = 0$$
 and  $pq \neq 0$  we have

$$\begin{split} \Phi_{25}(\xi) &= \frac{-pd}{q[d+\cosh(p\xi)-\sinh(p\xi)]},\\ \Phi_{26}(\xi) &= -\frac{p[\cosh(p\xi)+\sinh(p\xi)]}{q[d+\cosh(p\xi)+\sinh(p\xi)]}, \end{split}$$

where *d* is an arbitrary constant.

**Type 4:** When r = p = 0 and  $q \neq 0$  we have

$$\Phi_{_{27}}(\xi) = \frac{-1}{q\xi + c_1} ,$$

where c<sub>1</sub> is an arbitrary constant.

**Step 6:** Substituting the well known solutions of Equation (6) listed above in Step 5 into (5) we have many families of exact solutions of Equation (2).

## MANY FAMILIES OF EXACT TRAVELING WAVE SOLUTIONS FOR EQUATION (1)

Here we apply the proposed improved generalized Riccati equation mapping method to find many families of exact traveling wave solutions of Equation (1). To the end we use the wave transformation

$$u(x,t) = u(\xi), \quad \xi = \frac{1}{L}x - \frac{c}{\tau}t,$$
 (7)

where  $\tau = R_1 C_0 = 1.32 \times 10^{-6} s$ , and c is the dimensionless velocity of the wave, to reduce Equation (1) into the following ODE:

$$u'' - \frac{\alpha}{c}u' + \frac{\beta}{2}u^2 - \gamma u = 0 \tag{8}$$

where  $\alpha = \frac{\tau}{R_2 C_0}$ ,  $\beta = \frac{2R_1\delta}{R_2}$ ,  $\gamma = \frac{R_1}{R_2}$ .

By balancing  $u''with u^2$ , we have m=2. Hence the formal solution of Equation (8) takes the form:

$$u(\xi) = a_2 Q^2 + a_1 Q + a_0 + a_{-1} Q^{-1} + a_{-2} Q^{-2}$$
(9)

where  $a_2$ ,  $a_1$ ,  $a_0$ ,  $a_{-1}$ ,  $a_{-2}$  are constants to be determined, such that  $a_{-2} \neq 0$  or  $a_2 \neq 0$ .

Inserting (9) with the aid of Equation (6) into Equation (8) we get the following system of algebraic equations:

$$Q^{4} : \qquad 6a_{2}q^{2} + \frac{\beta}{2}a_{2}^{2} = 0,$$
$$Q^{-4} : \qquad 6a_{-2}r^{2} + \frac{\beta}{2}a_{-2}^{2} = 0,$$

$$Q^{3} : 10a_{2}pq + 2a_{1}q^{2} - \frac{2q\alpha}{c}a_{2} + \beta a_{1}a_{2} = 0,$$

$$Q^{-3} : 10a_{-2}pr + 2a_{-1}r^{2} + \frac{2r\alpha}{c}a_{-2} + \beta a_{-1}a_{-2} = 0,$$

$$Q^{2} : 8a_{2}qr + 3a_{1}pq + 4a_{2}p^{2} - \frac{\alpha}{c}(2a_{2}p + a_{1}q) + \frac{\beta}{2}(a_{1}^{2} + 2a_{0}a_{2}) - \gamma a_{2} = 0,$$

$$Q^{-2} : 8a_{-2}qr + 3a_{-1}pr + 4a_{-2}p^{2} + \frac{\alpha}{c}(2a_{-2}p + a_{-1}r) + \frac{\beta}{2}(a_{-1}^{2} + 2a_{0}a_{-2}) - \gamma a_{-2} = 0,$$

$$Q : 6a_{2}pr + 2a_{-1}qr + a_{1}p^{2} - \frac{\alpha}{c}(2a_{2}r + a_{1}p) + \beta(a_{-1}a_{2} + a_{0}a_{1}) - \gamma a_{1} = 0,$$

$$Q^{-1} : 6a_{-2}pq + 2a_{-1}qr + a_{-1}p^{2} + \frac{\alpha}{c}(2a_{-2}q + a_{-1}p) + \beta(a_{1}a_{-2} + a_{0}a_{-1}) - \gamma a_{-1} = 0,$$

$$Q^{0} : 2a_{2}r^{2} + a_{1}pr + a_{-1}pq + 2a_{-2}q^{2} - \frac{\alpha}{c}(a_{1}r - a_{-1}q) + \frac{\beta}{2}(a_{0}^{2} + 2a_{-1}a_{1} + 2a_{-2}a_{2}) - \gamma a_{0} = 0.$$

By solving these algebraic equations with the aid of Maple or Mathematica we have the following cases:

#### Case 1

$$p = p, q = q, r = -\frac{1}{24q}(\gamma - 6p^2), a_0 = \frac{3}{2\beta}(-2p^2 + \frac{10}{3}\frac{p\gamma c}{\alpha} + \gamma), a_1 = \frac{12q(\frac{\alpha}{5c} - p)}{\beta},$$
$$a_2 = \frac{-12q^2}{\beta}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}, a_{-1} = a_{-2} = 0$$

#### Case 2

$$p = p, q = q, r = -\frac{1}{24q}(\gamma - 6p^2), a_0 = \frac{1}{2\beta}(-6p^2 - \frac{10p\gamma c}{\alpha} + 3\gamma), a_{-1} = \frac{(\frac{\alpha}{5c} + p)(\gamma - 6p^2)}{2q\beta}, a_{-2} = \frac{-(\gamma - 6p^2)^2}{48\beta q^2}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}, a_1 = a_2 = 0$$

#### Case 3

$$p = 0, q = q, r = -\frac{\gamma}{96q}, a_0 = \frac{5\gamma}{4\beta}, a_1 = \frac{12q\alpha}{5\beta c}, a_{-1} = \frac{\gamma\alpha}{40q\beta c}, a_2 = \frac{-12q^2}{\beta}$$
$$a_{-2} = \frac{-\gamma^2}{768\beta q^2}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}.$$

### Exact traveling wave solutions of Equation (1) for Case 1

By using the case 1 and according to the values of solutions of type 1 in the proposed method, we obtain the following exact traveling wave solutions for Equation (1):

$$u_{1}(x,t) = \frac{3\gamma}{2\beta} + 2p[\frac{5\gamma c}{2\alpha} - \frac{3\alpha}{5\beta c}] - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \tanh(\sqrt{\frac{\gamma}{24}}\xi) - \frac{\gamma}{2\beta} \tanh^{2}(\sqrt{\frac{\gamma}{24}}\xi),$$
$$u_{2}(x,t) = \frac{3\gamma}{2\beta} + 2p[\frac{5\gamma c}{2\alpha} - \frac{3\alpha}{5\beta c}] - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \coth(\sqrt{\frac{\gamma}{24}}\xi) - \frac{\gamma}{2\beta} \coth^{2}(\sqrt{\frac{\gamma}{24}}\xi),$$

$$\begin{split} u_{3}(x,t) &= \frac{3\gamma}{2\beta} + 2p[\frac{5\gamma c}{2\alpha} - \frac{3\alpha}{5\beta c}] - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \bigg[ \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \, \operatorname{sec} h(\sqrt{\frac{\gamma}{6}}\xi) \bigg] \\ &\quad - \frac{\gamma}{2\beta} \bigg[ \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \, \operatorname{sec} h(\sqrt{\frac{\gamma}{6}}\xi) \bigg]^{2}, \\ u_{4}(x,t) &= \frac{3\gamma}{2\beta} + 2p[\frac{5\gamma c}{2\alpha} - \frac{3\alpha}{5\beta c}] - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \bigg[ \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm i \, \operatorname{cos} ch(\sqrt{\frac{\gamma}{6}}\xi) \bigg] \\ &\quad - \frac{\gamma}{2\beta} \bigg[ \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm i \, \operatorname{cos} ch(\sqrt{\frac{\gamma}{6}}\xi) \bigg]^{2}, \\ u_{5}(x,t) &= \frac{3\gamma}{2\beta} + p[\frac{5\gamma c}{\alpha\beta} - \frac{6\alpha}{5\beta c}] - \frac{3\alpha}{5\beta c} \sqrt{\frac{\gamma}{6}} \bigg[ \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm i \, \operatorname{cos} ch(\sqrt{\frac{\gamma}{96}}\xi) \bigg] \\ &\quad - \frac{\gamma}{2\beta} \bigg[ \coth(\sqrt{\frac{\gamma}{96}}\xi) \pm i \, \operatorname{cos} ch(\sqrt{\frac{\gamma}{96}}\xi) \bigg]^{2}, \\ u_{5}(x,t) &= \frac{3\gamma}{2\beta} + p[\frac{5\gamma c}{\alpha\beta} - \frac{6\alpha}{5\beta c}] - \frac{3\alpha}{5\beta c} \sqrt{\frac{\gamma}{6}} \bigg[ \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \bigg] \\ &\quad - \frac{\gamma}{8\beta} \bigg[ \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \bigg]^{2}, \\ u_{6}(x,t) &= \frac{3\gamma}{2\beta} + \frac{p}{\beta} [\frac{5\gamma c}{\alpha} - \frac{6\alpha}{5c}] + \frac{\alpha\sqrt{6\gamma}}{5\beta c} \bigg[ \frac{(\sqrt{A^{2} + B^{2}} - A \cosh(\sqrt{\frac{\gamma}{6}}\xi))}{(A \sinh(\sqrt{\frac{\gamma}{6}}\xi) + B)} \bigg] - \frac{\gamma}{2\beta} \bigg[ \frac{(\sqrt{A^{2} + B^{2}} - A \cosh(\sqrt{\frac{\gamma}{6}}\xi))}{(A \sinh(\sqrt{\frac{\gamma}{6}}\xi) + B)} \bigg]^{2} \\ u_{7}(x,t) &= \frac{3\gamma}{2\beta} + \frac{p}{\beta} [\frac{5\gamma c}{\alpha} - \frac{6\alpha}{5c}] + \frac{\alpha\sqrt{6\gamma}}{5\beta c} \bigg[ \frac{(\sqrt{B^{2} - A^{2}} + A \cosh(\sqrt{\frac{\gamma}{6}}\xi))}{(A \sinh(\sqrt{\frac{\gamma}{6}}\xi) + B)} \bigg] - \frac{\gamma}{2\beta} \bigg[ \frac{(\sqrt{B^{2} - A^{2}} + A \cosh(\sqrt{\frac{\gamma}{6}}\xi))}{(A \sinh(\sqrt{\frac{\gamma}{6}}\xi) + B)} \bigg]^{2} \end{split}$$

where A and B are two non-zero real constants satisfying  $B^2-A^2>0$ ,

$$\begin{split} u_{8}(x,t) &= \frac{3}{2\beta} [-2p^{2} + \frac{10p\gamma c}{3\alpha} + \gamma] + \frac{24qr}{\beta} \left( \frac{(-p + \frac{\alpha}{5c})\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{(\sqrt{\frac{\gamma}{6}}\sinh(\sqrt{\frac{\gamma}{24}}\xi) - \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi))} \right) \\ &- \frac{48q^{2}r^{2}\cosh^{2}(\sqrt{\frac{\gamma}{24}}\xi)}{\beta \left(\sqrt{\frac{\gamma}{6}}\sinh(\sqrt{\frac{\gamma}{24}}\xi) - \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)\right)^{2}}, \\ u_{9}(x,t) &= \frac{3}{2\beta} [-2p^{2} + \frac{10p\gamma c}{3\alpha} + \gamma] - \frac{24qr}{\beta} \left( \frac{(-p + \frac{\alpha}{5c})\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{(p\sinh(\sqrt{\frac{\gamma}{24}}\xi) - \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi))} \right) \end{split}$$

$$-\frac{48q^2r^2\sinh^2(\sqrt{\frac{\gamma}{24}}\,\xi)-\sqrt{\frac{2}{6}}\cosh(\sqrt{\frac{2}{24}}\,\xi)}{\beta\left(p\sinh(\sqrt{\frac{\gamma}{24}}\,\xi)-\sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\,\xi)\right)^2},$$

$$\begin{split} u_{10}(x,t) &= \frac{3}{2\beta} [-2p^2 + \frac{10p\gamma\epsilon}{3\alpha} + \gamma] + \frac{24qr}{\beta} \Biggl( \frac{(-p + \frac{\alpha}{5\epsilon}) \cosh(\sqrt{\frac{\gamma}{24}}\xi)}{(\sqrt{\frac{\gamma}{6}} \sinh(\sqrt{\frac{\gamma}{6}}\xi) - p\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm i\sqrt{\frac{\gamma}{6}})} \Biggr) \\ &- \frac{48q^2r^2\cosh^2(\sqrt{\frac{\gamma}{24}}\xi)}{\beta \biggl(\sqrt{\frac{\gamma}{6}}\sinh(\sqrt{\frac{\gamma}{6}}\xi) - p\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm i\sqrt{\frac{\gamma}{6}})^2}, \end{split}$$

$$\begin{split} u_{11}(x,t) &= \frac{3}{2\beta} [-2p^2 + \frac{10p\,\gamma c}{3\alpha} + \gamma] + \frac{24qr}{\beta} \Biggl( \frac{(-p + \frac{\alpha}{5c})\sinh(\sqrt{\frac{\gamma}{24}}\,\xi)}{(-p\sinh(\sqrt{\frac{\gamma}{6}}\,\xi) - \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{6}}\,\xi) \pm \sqrt{\frac{\gamma}{6}})} \Biggr) \\ &- \frac{48q^2r^2\sinh^2(\sqrt{\frac{\gamma}{24}}\,\xi)}{\beta \biggl( -p\sinh(\sqrt{\frac{\gamma}{6}}\,\xi) - \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{6}}\,\xi) \pm \sqrt{\frac{\gamma}{6}}\,\beta \biggr)^2}, \end{split}$$
$$u_{12}(x,t) &= \frac{3}{2\beta} [-2p^2 + \frac{10p\,\gamma c}{3\alpha} + \gamma] + \frac{24qr}{\beta} \Biggl( \frac{(-p + \frac{\alpha}{5c})\sinh(\sqrt{\frac{\gamma}{24}}\,\xi)}{(-p\sinh(\sqrt{\frac{\gamma}{24}}\,\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\,\xi))} \Biggr) \\ &- \frac{48q^2r^2\sinh^2(\frac{\sqrt{\gamma}}{2\sqrt{6}}\,\xi)}{\beta \biggl( -p\sinh(\sqrt{\frac{\gamma}{24}}\,\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\,\xi) \biggr)^2}, \end{split}$$

where 
$$\xi = \frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau}$$
.

**Remark 1.** If  $\gamma = 6p^2$ , then r = 0 and  $\Delta = p^2$ . Consequently we have

$$a_0 = \frac{12P^2}{\beta}$$
,  $a_2 = \frac{-12q^2}{\beta}$ ,  $c = \frac{\alpha}{5p}$ ,  $a_{-1} = a_1 = a_{-2} = 0$ .

According to the values of the solutions of type 3 in Sec.2, we obtain the following exact traveling wave solutions for Equation (1):

$$u_{13}(x,t) = \frac{12p^2}{\beta} - \frac{12q^2}{\beta} \left( \frac{pd}{q[d + \cosh(p\xi) - \sinh(p\xi)]} \right)^2,$$
  
$$u_{14}(x,t) = \frac{12p^2}{\beta} - \frac{12q^2}{\beta} \left( \frac{p[\cosh(p\xi) + \sinh(p\xi)]}{q[d + \cosh(p\xi) - \sinh(p\xi)]} \right)^2,$$
  
where  $\xi = \frac{1}{L}x - \frac{\alpha}{5p}\frac{t}{\tau}.$ 

### Exact traveling wave solutions of Equation (1) for case 2.

By using the case 2 and according to the values of solutions of type 1 in the proposed method, we obtain the following exact traveling wave solutions for Equation (1):

$$u_{1}(x,t) = \frac{1}{2\beta} \left[ -6p^{2} - \frac{10p\gamma e}{\alpha} + 3\gamma \right] - \frac{(\gamma - 6p^{2})(p + \frac{\alpha}{5e})}{\beta} \left[ p + \sqrt{\frac{\gamma}{6}} \tanh(\sqrt{\frac{\gamma}{24}} \xi) \right]^{-1} - \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ p + \sqrt{\frac{\gamma}{6}} \tanh(\sqrt{\frac{\gamma}{24}} \xi) \right]^{-2}$$

$$\begin{split} u_{2}(x,t) &= \frac{1}{2\beta} [-6p^{2} - \frac{10p\gamma c}{\alpha} + 3\gamma] - \frac{(\gamma - 6p^{2})(p + \frac{\alpha}{5c})}{\beta} \left[ p + \sqrt{\frac{\gamma}{6}} \operatorname{coth}(\sqrt{\frac{\gamma}{24}} \xi) \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ p + \sqrt{\frac{\gamma}{6}} \operatorname{coth}(\sqrt{\frac{\gamma}{24}} \xi) \right]^{2} \\ u_{3}(x,t) &= \frac{1}{2\beta} [-6p^{2} - \frac{10p\gamma c}{\alpha} + 3\gamma] - \frac{(\gamma - 6p^{2})(p + \frac{\alpha}{5c})}{\beta} \left[ p + \sqrt{\frac{\gamma}{6}} [\tanh(\sqrt{\frac{\gamma}{6}} \xi) \pm i \operatorname{sec} h(\sqrt{\frac{\gamma}{6}} \xi) \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ p + \sqrt{\frac{\gamma}{6}} [\tanh(\sqrt{\frac{\gamma}{5}} \xi) \pm i \operatorname{sec} h(\sqrt{\frac{\gamma}{6}} \xi) \right]^{-2} \\ u_{4}(x,t) &= \frac{1}{2\beta} [-6p^{2} - \frac{10p\gamma c}{\alpha} + 3\gamma] - \frac{(\gamma - 6p^{2})(p + \frac{\alpha}{5c})}{\beta} \left[ p + \sqrt{\frac{\gamma}{6}} [\coth(\sqrt{\frac{\gamma}{6}} \xi) \pm \operatorname{csc} h(\sqrt{\frac{\gamma}{6}} \xi) \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ p + \sqrt{\frac{\gamma}{6}} [\coth(\sqrt{\frac{\gamma}{6}} \xi) \pm \operatorname{csc} h(\sqrt{\frac{\gamma}{6}} \xi) \right]^{-1} \\ u_{5}(x,t) &= \frac{1}{2\beta} [-6p^{2} - \frac{10p\gamma c}{\alpha} + 3\gamma] - \frac{2(\gamma - 6p^{2})(p + \frac{\alpha}{5c})}{\beta} \left[ 2p + \sqrt{\frac{\gamma}{96}} \xi [\tanh(\sqrt{\frac{\gamma}{96}} \xi) \pm \coth(\sqrt{\frac{\gamma}{96}} \xi) \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{3\beta} \left[ 2p + \sqrt{\frac{\gamma}{6}} [\tanh(\sqrt{\frac{\gamma}{96}} \xi) \pm \operatorname{coth}(\sqrt{\frac{\gamma}{96}} \xi) \right]^{-1} \\ u_{6}(x,t) &= \frac{1}{2\beta} [-6p^{2} - \frac{10p\gamma c}{\alpha} + 3\gamma] + \frac{(\gamma - 6p^{2})(p + \frac{\alpha}{5c})}{\beta} \left[ -p + \frac{\sqrt{\frac{\gamma}{6}} [\sqrt{A^{2} + B^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p + \frac{\sqrt{\frac{\gamma}{6}} [\sqrt{A^{2} + B^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p + \frac{\sqrt{\frac{\gamma}{6}} [\sqrt{A^{2} + B^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p - \frac{\sqrt{\frac{\gamma}{6}} [\sqrt{B^{2} - A^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p - \frac{\sqrt{\frac{\gamma}{6}} [\sqrt{B^{2} - A^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p - \sqrt{\frac{\sqrt{\frac{\gamma}{6}} [\sqrt{B^{2} - A^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p - \sqrt{\frac{\sqrt{\frac{\gamma}{6}} [\sqrt{B^{2} - A^{2}} - A \operatorname{cosh}(\sqrt{\frac{\gamma}{6}} \xi)]}}{A \operatorname{sinh}(\sqrt{\frac{\gamma}{6}} \xi) + B} \right]^{-1} \\ &- \frac{(\gamma - 6p^{2})^{2}}{12\beta} \left[ -p - \sqrt{\frac{\sqrt{\frac{\gamma}{6}$$

where A and B are two non-zero real constants satisfying  $B^2$ - $A^2$ >0,

$$\begin{split} u_{s}(x,t) &= \frac{1}{2\beta} \left[ -\delta p^{z} - \frac{10p\gamma c}{\alpha} + 3\gamma \right] - \frac{(\gamma - \delta p^{z})(p + \frac{\alpha}{5c})}{4q^{\gamma}\beta} \left[ \frac{\cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)}{\sqrt{\frac{\gamma}{6}} \sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right) - p \cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)} \right]^{-1} \\ &- \frac{(\gamma - \delta p^{z})^{z}}{192\beta q^{z\gamma z}} \left[ \frac{\cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)}{\sqrt{\frac{\beta}{6}} \sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right) - p \cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)} \right]^{-1} \\ u_{s}(x,t) &= \frac{1}{2\beta} \left[ -\delta p^{z} - \frac{10p\gamma c}{\alpha} + 3\gamma \right] - \frac{(\gamma - \delta p^{z})(p + \frac{\alpha}{5c})}{4q\gamma\beta} \left[ \frac{\sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)}{p \sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right) - \sqrt{\frac{\beta}{6}} \cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)} \right]^{-1} \\ &- \frac{(\gamma - \delta p^{z})^{z}}{192\beta q^{z\gamma z}} \left[ \frac{\sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)}{p \sinh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right) - \sqrt{\frac{\gamma}{6}} \cosh\left(\sqrt{\frac{\gamma}{24}}\,\xi\right)} \right]^{-1} \end{split}$$

$$\begin{split} u_{10}(x,t) &= \frac{1}{2\beta} \left[ -6p^2 - \frac{10p\gamma c}{\alpha} + 3\gamma \right] + \frac{(\gamma - 6p^2)(p + \frac{\alpha}{5c})}{4qr\beta} \left[ \frac{\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{\sqrt{\frac{\beta}{6}}\sinh(\sqrt{\frac{\gamma}{6}}\xi) - p\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm i\sqrt{\frac{\gamma}{6}}} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{\sqrt{\frac{\gamma}{6}}\sinh(\sqrt{\frac{\gamma}{6}}\xi) - p\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm i\sqrt{\frac{\gamma}{6}}} \right]^{-2} \\ u_{11}(x,t) &= \frac{1}{2\beta} \left[ -6p^2 - \frac{10p\gamma c}{\alpha} + 3\gamma \right] + \frac{(\gamma - 6p^2)(p + \frac{\alpha}{5c})}{4qr\beta} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{6}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm \sqrt{\frac{\gamma}{6}}} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{6}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{6}}\xi) \pm \sqrt{\frac{\gamma}{6}}} \right]^{-2} \\ u_{12}(x,t) &= \frac{1}{2\beta} \left[ -6p^2 - \frac{10p\gamma c}{\alpha} + 3\gamma \right] + \frac{(\gamma - 6p^2)(p + \frac{\alpha}{5c})}{4qr\beta} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-2} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}}\cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}} \cosh(\sqrt{\frac{\gamma}{24}}\xi)} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-p\sinh(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-phi(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-phi(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}} \right]^{-1} \\ &- \frac{(\gamma - 6p^2)^2}{192\beta q^2 r^2} \left[ \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{-phi(\sqrt{\frac{\gamma}{24}}\xi) + \sqrt{\frac{\gamma}{6}}} \right]^{-1} \\ &- \frac{(\gamma - 6p^2$$

where 
$$\xi = \frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{\gamma}{6}}\frac{t}{\tau}$$
.

# Exact traveling wave solutions of Equation (1) for Case 3

By using the case 3 and according to the values of solutions of type 1 in the proposed method, we obtain the following exact traveling wave solutions for Equation (1):

$$\begin{split} u_{1}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \tanh(\sqrt{\frac{\gamma}{24}}\xi) + \frac{12qr}{\beta} \tanh^{2}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\alpha\sqrt{6\gamma}}{20\betac} \coth(\sqrt{\frac{\gamma}{24}}\xi) - \frac{\gamma}{32\beta} \coth^{2}(\sqrt{\frac{\gamma}{24}}\xi) \\ u_{2}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \coth(\sqrt{\frac{\gamma}{24}}\xi) + \frac{12qr}{\beta} \coth^{2}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\alpha\sqrt{6\gamma}}{20\betac} \tanh(\sqrt{\frac{\gamma}{24}}\xi) - \frac{\gamma}{32\beta} \tanh^{2}(\sqrt{\frac{\gamma}{24}}\xi) \\ u_{3}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \left( \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \sec h(\sqrt{\frac{\gamma}{6}}\xi) \right) + \frac{12qr}{\beta} \left( \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \sec h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \sec h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-1} - \frac{\gamma}{32\beta} \left( \tanh(\sqrt{\frac{\gamma}{6}}\xi) \pm i \sec h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-2} \\ u_{4}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-1} - \frac{\gamma}{32\beta} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-1} + \frac{\gamma}{32\beta} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-1} + \frac{\gamma}{32\beta} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-1} + \frac{\gamma}{32\beta} \left( \coth(\sqrt{\frac{\gamma}{6}}\xi) \pm \csc h(\sqrt{\frac{\gamma}{6}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \sqrt{\frac{32}{2}} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} + \frac{\gamma}{8\beta} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} - \frac{\gamma}{8\beta} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} - \frac{\gamma}{8\beta} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} - \frac{\gamma}{8\beta} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} - \frac{\gamma}{8\beta} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm \coth(\sqrt{\frac{\gamma}{96}}\xi) \right)^{-1} \\ &+ \frac{\alpha\sqrt{6\gamma}}{10\betac} \left( \tanh(\sqrt{\frac{\gamma}{96}}\xi) \pm$$

$$\begin{split} u_{6}(x,t) &= \frac{5\gamma}{4\beta} + \frac{\alpha\sqrt{6\gamma}}{5\beta c} \left( \frac{\sqrt{A^{2} + B^{2}} - A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(2\sqrt{-qr}\xi) + B} \right) + \frac{12qr}{\beta} \left( \frac{\sqrt{A^{2} + B^{2}} - A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{2} \\ &+ \frac{\alpha\sqrt{6\gamma}}{20\beta c} \left( \frac{\sqrt{A^{2} + B^{2}} - A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{-1} - \frac{\gamma}{32\beta} \left( \frac{\sqrt{A^{2} + B^{2}} - A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{2} \\ u_{\gamma}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \left( \frac{\sqrt{B^{2} - A^{2}} + A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right) + \frac{12qr}{\beta} \left( \frac{\sqrt{B^{2} - A^{2}} + A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{2} \\ &- \frac{\alpha\sqrt{6\gamma}}{20\beta c} \left( \frac{\sqrt{B^{2} - A^{2}} + A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{-1} - \frac{\gamma}{32\beta} \left( \frac{\sqrt{B^{2} - A^{2}} + A\cosh(\sqrt{\frac{\gamma}{6}}\xi)}{A\sinh(\sqrt{\frac{\gamma}{6}}\xi) + B} \right)^{-2} \end{split}$$

where A and B are two non-zero real constants satisfying  $B^2-A^2>0$ ,

 $u_{s}(x,t) = \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \operatorname{coth}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{3q}{\beta r} \operatorname{coth}^{2}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\alpha\gamma}{20q\beta c}\sqrt{\frac{\gamma}{24}} \tanh(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\gamma^{2}r}{192\beta q} \tanh^{2}(\sqrt{\frac{\gamma}{24}}\xi)$  $u_{y}(x,t) = \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \tanh(\sqrt{\frac{\gamma}{24}}\xi) + \frac{3q}{\beta r} \tanh^{2}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\alpha\gamma}{20q\beta c}\sqrt{\frac{\gamma}{24}} \operatorname{coth}(\sqrt{\frac{\gamma}{24}}\xi) + \frac{\gamma^{2}r}{192\beta q} \operatorname{coth}^{2}(\sqrt{\frac{\gamma}{24}}\xi)$ 

$$\begin{split} u_{10}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \left( \frac{\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{\sinh(\sqrt{\frac{\gamma}{6}}\xi)\pm i} \right) + \frac{12qr}{\beta} \left( \frac{\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{\sinh(\sqrt{\frac{\gamma}{6}}\xi)\pm i} \right)^2 - \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \frac{\coth(\sqrt{\frac{\gamma}{24}}\xi)}{\sinh(\sqrt{\frac{\gamma}{6}}\xi)\pm i} \right)^1 \\ &- \frac{\gamma}{32\beta} \left( \frac{\cosh(\sqrt{\frac{\gamma}{24}}\xi)}{\sinh(\sqrt{\frac{\gamma}{6}}\xi)\pm i} \right)^2 \end{split}$$

$$\begin{split} u_{11}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\betac} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}\xi})}{\cosh(\sqrt{\frac{\gamma}{6}\xi}) \pm i} \right) + \frac{12qr}{\beta} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}\xi})}{\cosh(\sqrt{\frac{\gamma}{6}\xi}) \pm i} \right)^2 - \frac{\alpha\sqrt{6\gamma}}{20\betac} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}\xi})}{\cosh(\sqrt{\frac{\gamma}{6}\xi}) \pm i} \right)^{-1} \\ &- \frac{\gamma}{32\beta} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}\xi})}{\cosh(\sqrt{\frac{\gamma}{6}\xi}) \pm i} \right)^{-2} \end{split}$$

$$\begin{split} u_{12}(x,t) &= \frac{5\gamma}{4\beta} - \frac{\alpha\sqrt{6\gamma}}{5\beta c} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{\cosh(\sqrt{\frac{\gamma}{24}}\xi) - 1} \right) + \frac{12qr}{\beta} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{\cosh(\sqrt{\frac{\gamma}{24}}\xi) - 1} \right)^2 - \frac{\alpha\sqrt{6\gamma}}{20\beta c} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{\cosh(\sqrt{\frac{\gamma}{24}}\xi) - 1} \right)^{-1} - \frac{\gamma}{31\beta} \left( \frac{\sinh(\sqrt{\frac{\gamma}{24}}\xi)}{\cosh(\sqrt{\frac{\gamma}{24}}\xi) - 1} \right)^{-2} \end{split}$$

where,  $\xi = \frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{\gamma}{6}}\frac{t}{\tau}$ .

### Remark 2

We have noted in the two cases 1, 2 that  $\Delta = \frac{\gamma}{6} > 0$  while

in case 3 we have  $\Delta = \frac{\gamma}{24} > 0$ . This yields that all the solutions of type 2 in the proposed method cannot be

considered in this paper. Therefore, Equation (1) has no trigonometric function solutions.

### PHYSICAL EXPLANATIONS FOR SOME OBTAINED RESULTS

Here, we will present some graphs for the obtained solutions of Equation (1) by selecting some special values of the parameters in the exact solutions using the mathematical software Maple, which can be shown below in Figures 1 to 6. From these explicit solutions, we see that  $u_1(x,t)$  in both cases 1, 2 are kink shaped soliton solutions while  $u_2(x,t)$  in these two cases are singular kink shaped soliton solutions. The two solutions  $u_1(x,t)$ 

and  $u_2(x,t)$  in case 3 are kink-singular shaped soliton solutions. The graphical representations of these solutions are shown in the following figures:

#### **CONCLUSIONS AND COMMENTS**

Here we give some comments on the solutions (22) to (27) of the second model of microtubules (16) in Sekulic et al. (2011a). We have found that some of these solutions are incorrect. Thus, we will correct these solutions and then compare between some of our results in the present article and the corrected solutions as follows:

1) There is a minor error in Equation (18) of Sekulic et al. (2011a). The correction is to replace  $\beta$  by  $\frac{\beta}{2}$  in Equation

form:

2) The results of Case I in Sekulic et al. (2011a) do not satisfy the algebraic Equation (21). After a careful revision, we have shown that the correction of this case should be in the form:

$$b = -\frac{\gamma}{24}, a_0 = \frac{3\gamma}{4\beta}, a_1 = \frac{6\alpha}{5\beta c}, a_2 = \frac{-6}{\beta}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}, b_1 = b_2 = 0$$
 (10)

It is easy to see that the corrected values (10) satisfy the algebraic Equation (21) of Sekulic et al. (2011a).

Hence with replacing  $\beta$  by  $\frac{\beta}{2}$  in (10), the result (22) of Sekulic et al. (2011a) should be rewritten in the corrected

$$u(x,t) = \frac{3\gamma}{2\beta} - \frac{12\alpha}{5\beta\epsilon} \sqrt{\frac{\gamma}{24}} \tanh\left(\sqrt{\frac{\gamma}{24}} \left(\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{\gamma}{6}} \frac{t}{\tau}\right)\right) - \frac{\gamma}{2\beta} \tanh^2\left(\sqrt{\frac{\gamma}{24}} \left(\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{\gamma}{6}} \frac{t}{\tau}\right)\right)$$
(11)



Figure 1. The plot of  $u_1(x,t)$  of case 1, where  $p = 1, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86.$ 



Figure 2. The plot of  $u_2(x,t)$  of case 1, where  $p = 1, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86$ .

On comparing the corrected result (11) with our result  $u_1(x,t)$  of "Exact traveling wave solutions of Equation (1) for case 1", we conclude that they are equivalent if p = 0 and q = 1.

3) The results of Case II in Sekulic et al. (2011a) do not satisfy the algebraic Equation (21) too. After a careful revision, we have shown that the correction of this case should be in the form:

$$b = \frac{\gamma}{24}, a_0 = \frac{\gamma}{4\beta}, a_1 = \frac{6\alpha}{5\beta c}, a_2 = \frac{-6}{\beta}, c = \frac{\alpha}{5}\sqrt{\frac{-6}{\gamma}}, b_1 = b_2 = 0$$
 (12)



Figure 3. The plot of  $u_1(x,t)$  of case 2, where  $p = 1, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86.$ 



Figure 4. The plot of  $u_2(x,t)$  of case 2, where  $p = 1, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86.$ 

It is easy to see that the values of (12) satisfy the algebraic Equation (21) of Sekulic et al. (2011a). From the values of (12) we deduce that  $\gamma = \frac{-6\alpha^2}{25c^2}$  which is negative. This contradicts that  $\gamma = \frac{R_1}{R_2} > 0$ . Therefore, Case 2 in Sekulic et al. (2011a) should be rejected.

4) The results of Case III in Sekulic et al. (2011a) do not satisfy the algebraic Equation (21) too. After a careful revision, we have shown that the correction of this case should be in the form:



Figure 5. The plot of  $u_1(x,t)$  of case 3, where  $p = 0, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86.$ 



Figure 6. The plot of  $u_2(x,t)$  of case 3, where  $p = 0, q = 1, \alpha = 104.76, \beta = 142.86, \gamma = 142.86.$ 

$$b = -\frac{\gamma}{24}, a_0 = \frac{3\gamma}{4\beta}, \ b_1 = \frac{\alpha\gamma}{20\beta c}, b_2 = \frac{-\gamma^2}{96\beta}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}, \ a_1 = a_2 = 0$$
(13)

It is easy to see that the corrected values of (13) satisfy the algebraic Equation (21) of Sekulic et al. (2011a).

Hence with replacing  $\beta$  by  $\frac{\beta}{2}$  in (13), the result (24) of (Sekulic et al, 2011a) should be rewritten in the corrected form:

$$u(x,t) = \frac{3\gamma}{2\beta} - \frac{\alpha\gamma}{10\beta c} \sqrt{\frac{24}{\gamma}} \coth\left(\sqrt{\frac{\gamma}{24}} \left(\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau}\right)\right) - \frac{\gamma}{2\beta} \coth^2\left(\sqrt{\frac{\gamma}{24}} \left(\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau}\right)\right)$$
(14)

On comparing the corrected result (14) with our result  $u_1(x,t)$  of 'Exact traveling wave solutions of Equation (1) for case 2', we conclude that they are equivalent if p = 0 and q = 1.

5) The results of Case IV of Sekulic et al. (2011a) still do not satisfy the algebraic Equation (21). After a careful revision, we have shown that the correction of this case should be in the form:

$$b = \frac{\gamma}{24}, \ c = \frac{\alpha}{5} \sqrt{\frac{-6}{\gamma}}, \ a_1 = a_2 = 0, \ b_2 = \frac{-\gamma^2}{96\beta}, \ b_1 = -\frac{\alpha\gamma}{20\beta c}, \ a_0 = \frac{\gamma}{4\beta}$$
(15)

It is easy to see that the corrected values of (15) satisfy the algebraic Equation (21) of Sekulic et al. (2011a). From the values of Equation (15) we deduce that  $\gamma = \frac{-6\alpha^2}{25c^2}$ which is negative. This contradicts that  $\gamma = \frac{R_1}{R_2} > 0$ . Therefore, the case IV in Sekulic et al. (2011a) should be rejected.

6) The results of Case V in Sekulic et al. (2011a) also do not satisfy the algebraic Equation (21). After a careful revision, we have shown that the correction of this case should be

$$b = -\frac{\gamma}{96}, a_0 = \frac{5\gamma}{8\beta}, \ b_1 = \frac{\alpha\gamma}{80\beta c}, b_2 = \frac{-\gamma^2}{1536\beta}, c = \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}, \ a_1 = \frac{6\alpha}{5\beta c}, a_2 = -\frac{6}{\beta}$$
 16)

It is easy to see that the corrected values of (16) satisfy the algebraic Equation (21) of Sekulic et al. (2011a).

Hence with replacing  $\beta$  by  $\frac{\beta}{2}$  in Equation (13), the result (26) of Sekulic et al. (2011a) should be rewritten in the corrected form:

$$u(x,t) = \frac{5\gamma}{4\beta} - \frac{12\alpha}{5\beta c} \sqrt{\frac{\gamma}{96}} \tanh\left(\sqrt{\frac{\gamma}{96}} (\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau})\right) - \frac{\gamma}{8\beta} \tanh^2\left(\sqrt{\frac{\gamma}{96}} (\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau})\right) + \frac{\alpha\gamma}{40\beta c} \sqrt{\frac{96}{\gamma}} \coth\left(\sqrt{\frac{\gamma}{96}} (\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau})\right) - \frac{\gamma}{8\beta} \coth^2\left(\sqrt{\frac{\gamma}{96}} (\frac{1}{L}x - \frac{\alpha}{5}\sqrt{\frac{6}{\gamma}}\frac{t}{\tau})\right)$$

$$(17)$$

On comparing the corrected result (17) with our result  $u_1(x,t)$  of 'Exact traveling wave solutions of Equation (1) for Case 3', we conclude that they are equivalent if p = 0 and q = 1.

7) The results of case VI in Sekulic et al. (2011a) do not

also satisfy Equation (21). After a careful revision, we have shown that the correction of this case should be in the form:

$$b = \frac{\gamma}{96}, \ c = \frac{\alpha}{5} \sqrt{\frac{-6}{\gamma}}, \ b_2 = -\frac{\gamma^2}{1536\beta}, \ b_1 = -\frac{\alpha\gamma}{80\beta c}, \ a_2 = \frac{-6}{\beta}, \ a_1 = \frac{6\alpha}{5\beta c}, \ a_0 = \frac{3\gamma}{8\beta}$$
(18)

It is easy to see that the corrected values of (18) satisfy the algebraic Equation (21) of Sekulic et al. (2011a). From the values of Equation (18) we deduce that  $\gamma = \frac{-6\alpha^2}{25c^2}$ which is negative. This contradicts that  $\gamma = \frac{R_1}{R_2} > 0$ .

Therefore, the case VI in Sekulic et al. (2011) should be rejected.

From these discussions we deduce that our results in the present article are new and recover the well-known results obtained in Sekulic et al. (2011a) after its corrections obtained above.

#### **Conflict of Interests**

The author(s) have not declared any conflict of interests.

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#### REFERENCES

- Abdou MA (2007).The extended tanh method and its applications for solving nonlinear physical models. Appl. Math. Comput.190(1):988-996. http://dx.doi.org/10.1016/j.amc.2007.01.070
- Ablowitz MJ, Clarkson PA (1991). Solitons, nonlinear evolution equations and inverse scattering transform. Cambridge University Press, New York. http://dx.doi.org/10.1017/CBO9780511623998; PMid:9905818
- Aslan I (2010). A note on the (G'/G) -expansion method again. Appl. Math. Comput. 217(2):937-938. http://dx.doi.org/10.1016/j.amc.2010.05.097
- Aslan I (2011). Exact and explicit solutions to the discrete nonlinear Schrödinger equation with a saturable nonlinearity. Phys. Lett. A. 375(47):4214-4217. http://dx.doi.org/10.1016/j.physleta.2011.10.009
- Aslan I (2012a). Some exact solutions for Toda type lattice differential equations using the improved (G'/G)– expansion method. Math. Methods Appl. Sci. 35(4):474-481. http://dx.doi.org/10.1002/mma.1579
- Aslan I (2012b). The discrete (G'/G) -expansion method applied to the differential-difference Burgers equation and the relativistic Toda lattice system. Numer. Methods Par. Diff. Eqs. 28(1):127-137. http://dx.doi.org/10.1002/num.20611
- Aslan I (2011). Comment on application of exp-function method (3+1)dimensional nonlinear evolution equation.[Comput Math. Appl. 56(2008):1451-1456], Comput. Math. Appl. 61(6):1700-1703. http://dx.doi.org/10.1016/j.camwa.2011.01.043
- Ayhan B, Bekir A (2012). The (G'/G) -expansion method for the nonlinear lattice equations. Commun. Nonlin. Sci. Numer. Simul. 17(9):3490-3498. http://dx.doi.org/10.1016/j.cnsns.2012.01.009
- Bekir A (2008). Application of the (G'/G) -expansion method for nonlinear evolution equations. Phys. Lett. A. 372(19):3400-3406.

http://dx.doi.org/10.1016/j.physleta.2008.01.057

- Bekir A (2009). The exp-function for Ostrovsky equation. Int. J. Nonlin.Sci.Numer.Simul.10:735-739.http://dx.doi.org/10.1515/IJNSNS.2009.10.6.735
- Bekir A (2010). Application of exp-function method for nonlinear differential-difference equations. Appl. Math. Comput. 215(11):4049-4053. http://dx.doi.org/10.1016/j.amc.2009.12.003
- Chen Y, Wang Q (2005). Extended Jacobi elliptic function rational expansion method and abundant families of Jacobi elliptic function solutions to (1+1)-dimensional dispersive long wave equation. Chaos Solit. Fract. 24(3):745-757. http://dx.doi.org/10.1016/j.chaos.2004.09.014
- Fan E (2000). Extended tanh-function method and its applications to nonlinear equations. Phys. Lett. A. 277:212-218. http://dx.doi.org/10.1016/S0375-9601(00)00725-8
- Freedman H, Rezania V, Priel A, Carpenter E, Noskov SY, Tuszynski JA (2010). Model of Ionic Currents through Microtubule Nanopores and the Lumen. Phys. Rev. E. 81(5):051912. http://dx.doi.org/10.1103/PhysRevE.81.051912; PMid:20866266
- He JH, Wu XH.(2006). Exp-function method for nonlinear wave equations. Chaos Solit. Fract. 30(3):700-708. http://dx.doi.org/10.1016/j.chaos.2006.03.020
- Hirota R (1971). Exact solutions of the KdV equation for multiple collisions of solutions. Phys. Rev. Lett. 27:1192-1194. http://dx.doi.org/10.1103/PhysRevLett.27.1192
- Ilic DI, Sataric MV, Ralevic N (2009), Atomic and molecular physics: Microtubule as a transmission line for ionic currents, Chin. Phys. Lett. 26:073101-073103. http://dx.doi.org/10.1088/0256-307X/26/7/073101
- Kudryashov NA (1988). Exact solutions of a generalized evolution of wave dynamics. J. Appl. Math. Mech. 52:361-365. http://dx.doi.org/10.1016/0021-8928(88)90090-1
- Kudryashov NA (1990). Exact solutions of a generalized Kuramoto-Sivashinsky equation. Phys. Lett. A. 147(5-6):287-291. http://dx.doi.org/10.1016/0375-9601(90)90449-X
- Kudryashov NA (1991).On types of nonlinear non integrable equations with exact solutions. Phys. Lett. A. 155(4-5):269-275. http://dx.doi.org/10.1016/0375-9601(91)90481-M
- Liu S, Z. Fu Z, Zhao Q (2001). Jacobi elliptic function expansion method and periodic wave solutions of nonlinear wave equations. Phys. Lett. A. 289(1-2):69-74. http://dx.doi.org/10.1016/S0375-9601(01)00580-1
- Lu D (2005). Jacobi elliptic function solutions for two variant Boussinesq equations, Chaos Solit. Fract. 24(5): 1373. http://dx.doi.org/10.1016/j.chaos.2004.09.085
- Miura MR (1978). Backlund transformation, Springer, Berlin, Germany.
- Rogers C, Shadwick WF (1982). Baclund Transformation and Their Applications. Academic Press, New York. p. 161.
- Sataric MV, Ilic DI, Ralevic NM and Tuszynski JA (2009). A nonlinear model of ionic wave propagation along microtubules. Eur. Biophys. J. Biophys. Lett. 38:637-647. http://dx.doi.org/10.1007/s00249-009-0421-5; PMid:19259657
- Sataric MV, Sekulic DL, Zivanov MB (2010). Solitonic ionic currents along microtubules. J. Comput. Theor. Nanosci. 7:2281-2290. http://dx.doi.org/10.1166/jctn.2010.1609
- Sekulic DL, Sataric MV, Zivanov MB (2011a). Symbolic computation of some new nonlinear partial differential equations of nanobiosciences using modified extended tanh-function method. Appl. Math. Comput. 218:3499-3506. http://dx.doi.org/10.1016/j.amc.2011.08.096
- Sekulic DL, Sataric BM, Tuszynski JA, Sataric MV (2011b). Nonlinear ionic pulses along microtubules. Eur. Phys. J. E. Soft Matter 34(5):1-11. http://dx.doi.org/10.1140/epje/i2011-11049-0; PMid:21604102
- Sekulić DL, Satarić MV (2012). Microtubule as nanobioelectronic nonlinear circuit, Serbian J. Elect. Eng. 9:107-119. http://dx.doi.org/10.2298/SJEE1201107S
- Sekulic DL, Sataric MV,Zivanov MB, Bajic JS (2012). Soliton-like pulses along electrical nonlinear transmission line, Electr. Elect. Eng.121:53-58.
- Wang M, Li X, Zhang J (2008). The (G'/G)-expansion method and traveling wave solutions of nonlinear evolution equations in mathematical physics. Phys. Lett. A. 372:417-423. http://dx.doi.org/10.1016/j.physleta.2007.07.051
- Weiss J, Tabor T, Carnevale G (1983). The Painlevé property for partial

- differential equations. J. Math. Phys. 24(3):552-526. http://dx.doi.org/10.1063/1.525721
- Yusufoglu E (2008).New solitary for the MBBM equations using expfunction method. Phys. Lett. A. 372:442-446. http://dx.doi.org/10.1016/j.physleta.2007.07.062
- Yusufoglu E, Bekir A (2008). Exact solutions of coupled nonlinear Klein-Gordon equations. Math. Comput. Model. 48:1694-1700. http://dx.doi.org/10.1016/j.mcm.2008.02.007
- Zayed EME (2009). The (G<sup>'</sup>/G) -expansion method and its applications to some nonlinear evolution equations in the mathematical physics. J. Appl. Math. Comput. 30:89-103. http://dx.doi.org/10.1007/s12190-008-0159-8
- Zayed EME (2010). Traveling wave solutions for higher dimensional nonlinear evolution equations using the (G'/G)-expansion method. J. Appl. Math. Inf. 28:383-395.
- Zayed EME, Arnous AH (2013). Many exact solutions for nonlinear dynamics of DNA model using the generalized Riccati equation mapping method. Sci. Res. Essays 8:340-346.
- Zayed EME, Amer YA, Shohib RMA (2013). The improved Riccati equation mapping method for constructing many families of exact solutions for nonlinear partial differential equation on nanobioscience. Int. J. Phys. Sci. 8(22):1246-1255.

- Zhang S (2008). Application of exp-function method to highdimensional nonlinear evolution equation. Chaos Solit. Fract.38:270-276. http://dx.doi.org/10.1016/j.chaos.2006.11.014
- Zhang S, Xia T (2008). A further improved tanh function method exactly solving the (2+1)- dimensional dispersive long wave equations. Appl. Math. E-Notes 8:58-66.
- Zhang S, Tong JL, Wang W (2008). A generalized (G'/G)-Expansion method for the mKdV equation with variable coefficients. Phys. Lett. A. 372:2254-2257. http://dx.doi.org/10.1016/j.physleta.2007.11.026
- Zhu SD (2008). The generalized Riccati equation mapping method in nonlinear evolution equation: application to (2+1)-dimensional Boitilion-Pempinelle equation. Chaos Solit. Fract. 37:1335-1342. http://dx.doi.org/10.1016/j.chaos.2006.10.015